# The disc-like host galaxies of radio-loud narrow-line Seyfert 1s 

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#### Abstract

Until recently, relativistic jets were ubiquitously found to be launched from giant elliptical galaxies. However, the detection by the Fermi-LAT of $\gamma$-ray emission from radio-loud narrowline Seyfert 1 (RL-NLSy1) galaxies raised doubts on this relation. Here, we morphologically characterize a sample of 29 RL-NLSy1s (including $12 \gamma$-emitters, $\gamma$-NLSy1s) in order to find clues on the conditions needed by active galactic nuclei (AGNs) to produce relativistic jets. We use deep near-infrared images from the Nordic Optical Telescope and the ESO VLT to analyse the surface brightness distribution of the galaxies in the sample. We detected 72 per cent of the hosts ( 24 per cent classified as $\gamma$-NLSy1s). Although we cannot rule out that some RL-NLSy1s are hosted by dispersion-supported systems, our findings strongly indicate that RL-NLSy1 hosts are preferentially disc galaxies. 52 percent of the resolved hosts ( 77 percent non- $\gamma$-emitters and 20 per cent $\gamma$-emitters) show bars with morphological properties (long and weak) consistent with models that promote gas inflows, which might trigger nuclear activity. The extremely red bulges of the $\gamma$-NLSy1s, and features that suggest minor mergers in 75 per cent of their hosts, might hint to the necessary conditions for $\gamma$-rays to be produced. Among the features that suggest mergers in our sample, we find six galaxies that show offset stellar bulges with respect to their AGNs. When we plot the nuclear versus the bulge magnitude, RL-NLSy1s locate in the low-luminosity end of flat spectrum radio quasars, suggesting a similar accretion mode between these two AGN types.


Key words: galaxies: active - galaxies: bulges - galaxies: Seyfert-galaxies: structure-gamma-rays: galaxies.

## 1 INTRODUCTION

Because of the tight empirical relations observed between the black hole mass and different properties of its host galaxy bulge (Magorrian et al. 1998; Ferrarese \& Merritt 2000; Gebhardt et al. 2000; Tremaine et al. 2002; Gültekin et al. 2009), it is now widely accepted that there is a strong connection between the supermassive black holes (SMBHs) and their host galaxies. The link between the black hole and its host galaxy is thought to be the active galactic nucleus (AGN) activity, through feedback processes (either positive or negative; for a review see Fabian 2012; Heckman \& Best 2014). If this is the case, then we can assume that the more powerful the AGN, the stronger the influence on its host galaxy. In fact, a study by Olguín-Iglesias et al. (2016) (performed on $0.3<z<1.0$, strongly beamed AGNs, whose relativistic jets point towards the Earth, i.e. blazars) suggests that the AGNs impact their hosts (either by suppression or by triggering of star formation) in a magnitude that is proportional to the jet power.

[^0]Until recently, powerful relativistic radio jets were virtually only found to be hosted in elliptical galaxies (e.g. Stickel et al. 1991; Kotilainen, Falomo \& Scarpa 1998a,b; Scarpa et al. 2000; Urry et al. 2000; Kotilainen, Hyvönen \& Falomo 2005; Hyvönen et al. 2007; Olguín-Iglesias et al. 2016), which helped develop ideas on how jets, SMBHs, and their host galaxies evolve. However, recently, a few studies report on blazar-like disc hosts, that is to say, with fully developed relativistic jets, capable of emitting $\gamma$-ray photons (León Tavares et al. 2014; Kotilainen et al. 2016; OlguínIglesias et al. 2017). These blazar-like disc galaxies constitute a peculiar type of AGNs known as narrow-line Seyfert 1s (NLSy1s), characterized by narrower Balmer line full width at half-maximum (FWHM $(\mathrm{H} \beta)<2000 \mathrm{~km} \mathrm{~s}^{-1}$ ) than in normal Seyferts, flux ratios [ $\mathrm{O}_{\mathrm{III}}$ ]/H $\beta<3$, strong optical $\mathrm{Fe}_{\text {II }}$ lines (Fe II bump), and a soft X-ray excess (Osterbrock \& Pogge 1985; Orban de Xivry et al. 2011). Based on the FWHMs of their broad-line region lines and the continuum luminosity (Kaspi et al. 2000), their central black holes masses $\left(M_{\mathrm{BH}}\right)$ are estimated to range from $\sim 10^{6}$ to $\sim 10^{7} \mathrm{M}_{\odot}$ (Mathur et al. 2012); thus, their accretion rates are thought to be close to the Eddington limit. A small fraction has been found to
be radio loud (RL, 7 per cent, Komossa et al. 2006; Berton et al. 2019), and among them, a smaller fraction (so far, 15 galaxies) are $\gamma$-ray emitting NLSy1s (hereafter, $\gamma$-NLSy1s). RL- and $\gamma$-NLSy1s have also shown intra-night optical variability (e.g. Ojha, Krishna \& Chand 2019), a tracer of jet activity.

RL-NLSy1s (including $\gamma$-NLSy1s) are excellent laboratories to study the mechanisms that make AGNs able to launch and collimate fully developed relativistic outflows at a likely early evolutionary stage. Thus, in this study, we characterize a sample of RL-NLSy1s (including $12 \gamma$-NLSy1s detected so far) with the aim of determining the properties of their host galaxies that could shed some light on the necessary conditions and mechanisms to generate the relativistic jet phenomenon.

This paper is organized as follows. In Section 2, we present the sample and observations. In Section 3, we explain the data reduction and the methodology of the analysis. In Section 4, we discuss the results and compare them with previous studies. Finally, in Section 6, we summarize our findings. All quantitative values given in this paper are based on a cosmology with $H_{0}=70 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{Mpc}^{-1}$, $\Omega_{\mathrm{m}}=0.3$, and $\Omega_{\Lambda}=0.7$.

## 2 SAMPLE AND OBSERVATIONS

The initial sample consists of $12 \gamma$-ray emitting NLSy1s (Abdo et al. 2009a,b; Foschini 2011; D'Ammando et al. 2012, 2015; Liao et al. 2015; Yao et al. 2015; Paliya et al. 2018). Given that $\gamma$ NLSy1s are also RL, we expanded this sample by imaging the host galaxies of 17 RL , but not $\gamma$-ray emitting, NLSy1s as a comparison sample. These galaxies are all observable from the Northern hemisphere, and have redshifts $z<0.5$ and radio-loudness ( $\mathrm{RL} \equiv f_{v, 4.85 \mathrm{GHz}} / f_{v, \mathrm{~B}}, \mathrm{RL}>31$ ).

The observations were conducted using two different telescopes, the Nordic Optical Telescope (NOT), using the near-infrared camera NOTcam [pixel scale $=0.234$ arcsec pixel ${ }^{-1}$ and field of view $($ FOV $)=4 \operatorname{arcmin} \times 4$ arcmin] and the ESO very large telescope (VLT), using its infrared spectrometer and array camera (ISAAC; pixel scale $=0.148$ arcsec and FOV $=152 \operatorname{arcsec} \times 152$ arcsec, Moorwood et al. 1998), depending on the declination of each galaxy. On the other hand, the RL-NLSy1s (but not $\gamma$-ray emitters) in the sample were observed with the NOT telescope, between 2013 January 23 and 2016 February 14 using the NOTcam.

As is usual in the near-infrared, all the targets in the sample were imaged using a jitter procedure to obtain a set of offset frames with respect to the initial position. Each target was observed in $J$ and/or $K$ bands, during an average exposure time of EXPTIME $\approx 1800 \mathrm{~s}$, and an average seeing $\sim 0.75 \operatorname{arcsec}$ (see Table 1).

## 3 DATA REDUCTION AND ANALYSIS

### 3.1 Data reduction

The images of the galaxies in the sample observed with the NOT were reduced using the NOTCam script for iraf reduce. ${ }^{1}$ This script takes the consecutive dithered images, corrects for flatfield, interpolates over bad pixels, and makes a sky template that is subtracted from each image. The images are then registered based on interactively selected stars (and RA/Dec. header keywords) and combined to obtain the final reduced image. For the images observed with the ISAAC, a similar procedure was followed. A flat frame was

[^1]derived from the twilight images and a sky image was obtained by median filtering the individual frames in the stack. The individual frames were then aligned using bright stars as reference points in the field and combined to produce the final reduced co-added image (see Olguín-Iglesias et al. 2016, for more details). The photometric calibration was performed by using the field stars in our images with the magnitudes reported by 2MASS (Skrutskie et al. 2006).

### 3.2 Photometric decomposition

The 2D light distribution of the reduced images is modelled using the image analysis algorithm Galfit (Peng et al. 2011). The different components of a galaxy are described using different analytical functions. For bulges and possible bars in the host galaxies of the sample, we use the Sérsic profile, whose functional form is
$\Sigma(r)=\Sigma_{\mathrm{e}} \exp \left[-\kappa\left(\left(\frac{r}{r_{\mathrm{e}}}\right)^{1 / n}-1\right)\right]$,
where $\Sigma_{\mathrm{e}}$ is the surface brightness of the pixel in the effective radius ( $r_{\mathrm{e}}$, radius where half of the total flux is concentrated). The parameter $n$ (the Sérsic index) is often referred to as a concentration parameter and the variable $\kappa$ is coupled to it.

We also use the exponential function, since it describes well the radial behaviour of galactic discs. Although the exponential function is a special case of the Sérsic profile (when $n=1$ ), nomenclaturewise, we use it when the component to fit is a disc, otherwise we use the Sérsic profile with $n=1$. Its functional form is
$\Sigma(r)=\Sigma_{0} \exp \left(-\frac{r}{r_{\mathrm{s}}}\right)$,
where $\Sigma(r)$ is the surface brightness at a radius $r, \Sigma_{0}$ is the surface brightness at the centre of the target, and $r_{\mathrm{s}}$ is the scale length of the disc.

The sky background is also modelled. We use a simple flat plane that can be tilted in $x$ and $y$ directions. Finally, the nuclear emission due to the powerful AGNs of the galaxies of the sample is fitted using a modelled point spread function (PSF).

### 3.2.1 PSF modelling

In most cases, the PSF modelling only consists of subtracting a bright ${ }^{2}$ non-saturated star close to the target and removing its background. However, it is not always possible to get a suitable star in the FOV. In the case where only saturated or faint stars are found, the following procedure is implemented:

First, we identify the stars in the FOV. Then, we select, preferentially, the stars with no sources within $\sim 7$ arcsec radius and more than $\sim 10$ arcsec away from the border of the FOV. The selected stars are centred in $50 \operatorname{arcsec} \times 50 \operatorname{arcsec}$ boxes, where all extra sources are masked out by means of the segmentation image process of SEXTRACTOR (Bertin \& Arnouts 1996). The wings of the PSF are modelled by fitting a saturated star with a number of exponential and Gaussian functions (Fig. 1, top panel). The core of the PSF is modelled by using another star (in this case, it is important not to be saturated) with Gaussian functions and the previously generated

[^2]Table 1. Main properties of the RL-NLSy1s analysed in this work and observations log.

| Source (1) | Name (2) | (3) | RA (4) | Dec. (5) | UT date (6) | Scale kpcarcmin ${ }^{-1}$ <br> (7) | Seeing (arcsec) <br> (8) | Exposure time <br> (s) <br> (9) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $0321+340^{a}$ | 1H 0323+342 | 0.061 | 03:24:41.1 | +34:10:46.0 | 23-Jan-13 | 1.177 | 1.00/1.00 | 1800/1800 |
| $0846+513^{a}$ | SBS 0846+513 | 0.580 | 08:49:58.0 | +51:08:29.0 | 14-Feb-16 | 6.579 | 0.65/0.67 | 3000/2700 |
| $0929+533^{a}$ | J093241+53063 | 0.590 | 09:32:41.1 | +53:06:33.3 | 30-Mar-18 | 6.633 | - /0.77 | -/4356 |
| $0948+002^{\text {a }}$ | PMN J0948+0022 | 0.585 | 09:48:57.3 | +00:22:26.0 | 21-Feb-14 | 6.606 | 0.73/0.77 | 4950/4890 |
| $0955+32^{\text {a }}$ | J095820+322401 | 0.530 | 09:58:20.9 | +32:24:01.6 | 31-Mar-18 | 6.293 | - /0.78 | -/3600 |
| 1102+2239 | FBQS J1102+2239 | 0.453 | 11:02:23.4 | +22:39:20.7 | 14-Feb-14 | 5.781 | -/0.79 | -/2700 |
| 1159-011 | IRAS 11598-0112 | 0.150 | 12:02:26.8 | - 01:29:15.0 | 13-Feb-16 | 2.614 | 0.79/0.82 | 2280/900 |
| 1200-004 | RXSJ12002-0046 | 0.179 | 12:00:14.1 | - 00:46:39.0 | 14-Feb-16 | 3.021 | 0.64/0.67 | 1080/960 |
| 1217+654 | J12176+6546 | 0.307 | 12:17:40.4 | +65:46:50.0 | 13-Feb-16 | 4.526 | 0.94/0.79 | 2160/2160 |
| 1219-044 ${ }^{\text {a }}$ | 4C+04.42 | 0.996 | 12:22:22.5 | +04:13:16.0 | 29-Mar-14 | 8.001 | 0.80/0.84 | 3650/2160 |
| $1227+321$ | RXSJ12278+3215 | 0.137 | 12:27:49.2 | +32:14:59.0 | 14-Feb-16 | 2.423 | 0.70/0.67 | 930/930 |
| $1246+0238^{\text {a }}$ | SDSS J124634.65+023809 | 0.362 | 12:46:34.6 | 02:38:09.1 | 12-Aug-17 | 5.048 | - /0.80 | -/3480 |
| $1337+600$ | J13374+6005 | 0.234 | 13:37:24.4 | +60:05:41.0 | 13-Feb-16 | 3.722 | 0.73/0.72 | 2160/2160 |
| $1403+022$ | J14033+0222 | 0.250 | 14:03:22.1 | +02:22:33.0 | 14-Feb-16 | 3.910 | 0.66/0.59 | 1380/1080 |
| $1421+3855^{a}$ | J142106+385522 ${ }^{\text {a }}$ | 0.490 | 14:21:06.0 | +38:55:22.5 | 14-Mar-18 | 6.038 | -/0.75 | -/2754 |
| 1441-476 ${ }^{\text {a }}$ | B3 $1441+476$ | 0.705 | 14:43:18.5 | +47:25:57.0 | 24-Apr-16 | 7.166 | - /0.80 | -/2700 |
| $1450+591$ | J14506+5919 | 0.202 | 14:50:42.0 | +59:19:37 | 22-May-14 | 3.325 | 0.82/0.82 | 900/1020 |
| $1502+036^{\text {a }}$ b | PKS 1502+036 | 0.409 | 15:05:06.5 | +03:26:31.0 | 04-Apr-13 | 5.446 | 0.80/0.82 | 920/280 |
| $1517+520$ | SBS 1517+520 | 0.371 | 15:18:32.9 | +51:54:57.0 | 14-Feb-16 | 5.128 | 0.63/0.67 | 1140/960 |
| $1546+353$ | B2 1546+35A | 0.479 | 15:48:17.9 | +35:11:28.0 | 13-Feb-16 | 5.964 | 0.68/0.74 | 2160/2340 |
| $1629+400$ | J16290+4007 | 0.272 | 16:29:01.3 | +40:08:00.0 | 13-Feb-16 | 4.157 | 0.63/0.62 | 2160/2160 |
| $1633+471$ | RXSJ16333+4718 | 0.116 | 16:33:23.5 | +47:19:00.0 | 14-Feb-16 | 2.101 | 0.62/0.63 | 1260/1050 |
| $1640+534$ | 2E 1640+5345 | 0.140 | 16:42:00.6 | +53:39:51.0 | 14-Feb-16 | 2.468 | 0.60/0.64 | 1200/1080 |
| $1641+345$ | J16410+3454 | 0.164 | 16:41:00.1 | +34:54:52.0 | 14-Feb-16 | 2.814 | 0.62/0.70 | 1320/1800 |
| $1644+261^{a}$ | FBQS J1644+2619 | 0.145 | 16:44:42.5 | +26:19:13.0 | 01-May-15 | 2.541 | 0.75/0.63 | 2550/2160 |
| $1702+457$ | B31702+457 | 0.060 | 17:03:30.3 | +45:40:47.2 | 21-Jun-16 | 1.159 | - /0.79 | -/945 |
| $1722+565$ | J17221+5654 | 0.426 | 17:22:06.0 | +56:54:51.0 | 13-Feb-16 | 5.579 | 0.62/0.60 | 2040/2040 |
| 2004-447 ${ }^{\text {a,b }}$ | PKS 2004-447 | 0.240 | 20:07:55.2 | - 44:34:44.0 | 16-Apr-13 | 3.793 | 0.40/0.45 | 600/110 |
| 2245-174 | IRAS 22453-1744 | 0.117 | 22:48:04.2 | - 17:28:30.0 | 14-Feb-16 | 2.116 | 1.44/- | 1260/- |

Notes. Columns: (1) and (2) give the designation and name of the source; (3) the redshift of the object; (4) and (5) the J2000 right ascension and declination of the source; (6) the observation date; (7) the target scale; (8) the seeing during the observation in $J$ and $K$ s band, respectively, and (9) the total exposure time for $J$ and $K \mathrm{~s}$ band, respectively.
${ }^{a} \gamma$-ray emitting NLSy1 galaxies.
${ }^{b}$ Galaxies observed using the ISAAC on the ESO/VLT.
wings model (Fig. 1, bottom panel). The magnitude difference between the saturated and non-saturated stars is important, since there must be an overlap in order to match the wings and core models. The resultant model is tested by fitting random stars in the FOV.

In order to take into account the image PSF in the modelling of the galaxies, we convolved our PSF model with the analytical functions used in the fitting. The PSF model is also used to fit the nuclear component, which, in the galaxies of this sample, is composed by the AGNs.

### 3.2.2 Uncertainties

The uncertainties of the real functional form of a given galaxy component lead to the errors in the galaxy models derived using the above-mentioned method. In order to estimate these errors, we follow the procedure described by Olguín-Iglesias et al. (2016), Kotilainen et al. (2016), and Olguín-Iglesias et al. (2017) who identify three sources of uncertainties: the PSF model, the sky background, and the zero-point of photometric calibration.
The uncertainties due to the PSF are accounted for by using different PSF models. A number of PSF models can be derived using different stars or different amount of Gaussians and exponentials. The uncertainties due to the sky is derived by running several sky fits in separated regions of 300 pixels $\times 300$ pixels in the FOV. The zero-
point magnitude depends on the star used to derive it since they are estimated from the magnitudes of several stars (see Section 3.1), then we use the zero-point magnitude variations ( $\sim \pm 0.1 \mathrm{mag}$ ) as another source of uncertainties.

Using these variations, GALFIT is run several times. The resultant fits are used to make a statistic where the best-fitting value is the mean and the errors are $\pm 1 \sigma$ for every parameter of the galaxy model.

### 3.3 Simulations

In order to assure the suitability of the images to resolve galaxy bulges, we performed a set of simulations (Table A1). The simulated galaxies have a nuclear component, represented by a star within the FOV. They also have a Sérsic function, with $n=1$, which represents a bulge and an exponential function representing a disc. The background is taken from the same image as the star of the nuclear component. Every galaxy in the simulation has a different combination of parameters (bulge, disc, and nuclear signal to noise ratios; bulge effective radius and seeing). The simulated galaxies are modelled in an identical way to the real galaxies in our sample in order to find whether the quality in our images is good enough to allow an acceptable subtraction of parameters. We found that the ability to properly retrieve the parameters of the bulges depends


Figure 1. PSF modelling procedure. In the top panel, we show the PSF wings modelling by fitting Gaussians and exponential functions to the surface brightness of a saturated star (star 1). Since this star is saturated, there is no information on the PSF core. In the bottom panel, we show the PSF core modelling by fitting a Gaussian function to the surface brightness of a non-saturated star (star 2). As can be seen, one part of the wings in this star is under the detection limit (red horizontal segmented line) and another part is noisy, hence the need of the saturated star in order to derive the PSF wings.
both on the size of the bulge, with respect to the seeing, and on the brightness of the bulge with respect to the brightness of the other components. In this way, the parameters of faint bulges (with respect to the nucleus) might be properly retrieved provided that they are large enough (e.g. simulation 42). We note that, although the parameters of a bulge might not be properly retrieved, it might still be detected, although with not enough quality to retrieve its parameters. This means that in some simulations, we were able to properly characterize the nucleus and disc and also detect a residual (unable to be fitted) that we knew, beforehand, was the bulge (e.g. simulation 1,6 , and 57 ). We focus on the parameters of bulge and nucleus, although all our simulations include a disc, with the intention of studying the effect of this component on the fittings. Most of the time the disc is not bright enough to hamper the bulge fit. However, the brightness of the bulge and the nucleus can often be high enough to make the modelling and even detection of the disc difficult (e.g. simulations 29, 30, and 75).

In Fig. 2, we show the set of simulations in a plot of the ratio between the bulge and nucleus $\mathrm{S} / \mathrm{N}$ versus bulge effective radius normalized to the seeing FWHM. We use signal-to-noise ratios


Figure 2. Ratio of the bulge over nucleus $\mathrm{S} / \mathrm{N}$ versus bulge effective radius normalized to the seeing FWHM. We see the location of the correctly retrieved simulated galaxies (green circles), the incorrectly retrieved simulated galaxies (red crosses), and the real galaxies in our sample (yellow stars, including all 18 galaxies with $K$-band observations from the NOTcam and fitted in this study, and excluding four galaxies from previous studies and/or observed with the ESO/ISAAC, six unresolved galaxies, and one observed only in $J$ band) A limit that divides the correctly from incorrectly retrieved simulated galaxies is represented by a segmented line.
instead of magnitudes because in our simulations we take into account the seeing. The blurring caused by seeing makes the signal of point sources smaller. It also spreads out the light from extended sources, which reduces the measured signal-to-noise and the ability of telescopes to see detail but does not affect the derived magnitudes. Hence, the same source observed with different seeings have the same magnitude, but different signal-to-noise ratios, which is crucial in galaxy fitting. In our simulations, signal-to-noise ratios refer to the peak counts of the source divided by the background noise in the same region. We find the bulk of our sample in the region where the correctly retrieved simulated galaxies lie, which gives us confidence in the reliability of our analysis. We consider a model correct when the difference between a given parameter of the simulation and the same parameter of the model is less than the maximum error for such parameter in our analysis of the real galaxies sample (i.e. magnitude error $\mathrm{Mag}_{\text {err }} \leq 0.50$, bulge effective radius error $\mathrm{Re}_{\text {err }} \leq 100$ per cent, and Sérsic index error $n_{\text {err }} \leq 1.35$ ). Two galaxies in our sample lie outside that region (PMNJ0948+002 and J095820+322401), suggesting that those parameters are incorrectly retrieved, and hence left out of the analysis.

## 4 THE HOST GALAXIES

In Table 2, we show a summary of the NLSy1 host and nuclear properties derived from the two-dimensional surface brightness decomposition analysis. In the cases where the host galaxy is not detected we estimated upper limits for their magnitudes by simulating a galaxy (following the method from Kotilainen et al. 2007; Olguín-Iglesias et al. 2016), assuming an average effective radius and Sérsic index from the successfully detected hosts, and then increasing the magnitude until the component becomes detectable with a signal-to noise ratio $\mathrm{S} / \mathrm{N}=5$, which by our own experience, is required in order to properly retrieve the structural parameters of the galaxy components.

| Name <br> (1) | Band <br> (2) | $M_{\text {nuc }}$ <br> (3) | $M_{\text {bulge }}$ <br> (4) | $M_{\text {disc }}$ <br> (5) | $\operatorname{Re}(\mathrm{kpc})$ <br> (6) | $\begin{aligned} & \mu_{\mathrm{e}} \\ & (7) \end{aligned}$ | $\begin{gathered} n \\ (8) \end{gathered}$ | (9) | $f_{\text {bar }}$ <br> (10) | $r_{\text {bar }}(\mathrm{kpc})$ <br> (11) | $\begin{aligned} & \chi_{\text {best }} \\ & (12) \end{aligned}$ | $\begin{gathered} \chi_{\text {best }} / \chi_{\text {psf }} \\ (13) \end{gathered}$ | Int <br> (14) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $1 \mathrm{H} 0323+342^{\text {a }}$ | J/K | $\begin{gathered} -23.00 /-24.74 \\ (0.01 / 0.01) \end{gathered}$ | $\begin{gathered} -22.53 /-24.15 \\ (0.01 / 0.01) \end{gathered}$ | $\begin{gathered} -23.49 /-24.93 \\ (0.01 / 0.01) \end{gathered}$ | $\begin{gathered} 1.48 / 1.40 \\ (0.01 / 0.01) \end{gathered}$ | $\begin{aligned} & 16.89 / 15.14 \\ & (0.01 / 0.01) \end{aligned}$ | $\begin{gathered} 0.88 / 1.24 \\ (0.002) \end{gathered}$ | 3.36/3.15 | - | - | 1.280 | - | Y |
| SBS0846+513 ${ }^{\text {a }}$ | J/K | $\begin{gathered} -27.14 /-28.99 \\ (0.36 / 0.38) \end{gathered}$ | >-23.89/>-25.68 | - | - | - | - | - | - | - | $\begin{gathered} 1.646 / 2.296 \\ (0.001 / 0.002) \end{gathered}$ | 1 | Y |
| $\mathrm{J} 093241+53063^{a}$ | K | $\begin{gathered} -26.90 \\ (0.30) \end{gathered}$ | $\begin{gathered} -25.25 \\ (0.39) \end{gathered}$ | - | $\begin{gathered} 5.57 \\ (0.58) \end{gathered}$ | $\begin{aligned} & 20.50 \\ & (0.48) \end{aligned}$ | $\begin{gathered} 1.21 \\ (0.70) \end{gathered}$ | - | - | - | $\begin{gathered} 1.167 \\ (0.001) \end{gathered}$ | 0.40 | Y |
| PMNJ0948 $+0022^{a, b}$ | J/K | $\begin{gathered} -27.39 /-29.50 \\ (0.40 / 0.39) \end{gathered}$ | $\begin{gathered} -24.53 /-26.91 \\ (0.39 / 0.41) \end{gathered}$ | $\begin{gathered} -25.19 /-27.59 \\ (0.40 / 0.41) \end{gathered}$ | $\begin{gathered} 1.82 / 0.90 \\ (0.33 / 0.36) \end{gathered}$ | $\begin{gathered} 19.5 / 17.5 \\ (0.33 / 0.30) \end{gathered}$ | $\begin{gathered} 1.00 / 1.50 \\ (0.28 / 0.28) \end{gathered}$ | $\begin{aligned} & 10.04 / 18.66 \\ & (0.36 / 0.38) \end{aligned}$ | - | - | $\begin{gathered} 1.718 \\ (0.002) \end{gathered}$ | 0.36 | N |
| $\mathrm{J} 095820+322401^{a, b}$ | K | $\begin{gathered} -24.11 \\ (0.30) \end{gathered}$ | $\begin{gathered} -22.91 \\ (0.36) \end{gathered}$ | $\begin{gathered} -25.94 \\ (0.48) \end{gathered}$ | $\begin{gathered} 1.47 \\ (0.41) \end{gathered}$ | $\begin{aligned} & 16.76 \\ & (0.50) \end{aligned}$ | $\begin{gathered} 0.60 \\ (0.89) \end{gathered}$ | $\begin{aligned} & 10.04 \\ & (0.36) \end{aligned}$ | - | - | $\begin{gathered} 1.200 \\ (0.001) \end{gathered}$ | 0.35 | Y |
| FBQS J1102+2239 | J/K | $\begin{gathered} -25.09 /-26.78 \\ (0.28 / 0.29) \end{gathered}$ | $\begin{gathered} -23.78 /-25.41 \\ (0.32 / 0.30) \end{gathered}$ | $\begin{gathered} -23.53 /-25.74 \\ (0.31 / 0.34) \end{gathered}$ | $\begin{gathered} 4.08 / 4.78 \\ (0.42 / 0.40) \end{gathered}$ | $\begin{gathered} 21.5 / 19.6 \\ (0.33 / 0.31) \end{gathered}$ | $\begin{gathered} 1.00 / 1.00 \\ (0.31 / 0.31) \end{gathered}$ | $\begin{aligned} & 11.06 / 16.24 \\ & (0.41 / 0.40) \end{aligned}$ | - | - | $\begin{gathered} 1.200 \\ (0.002) \end{gathered}$ | 0.32 | Y |
| IRAS11598-0112 | J/K | $\begin{gathered} -23.43 /-25.69 \\ (0.39 / 0.37) \end{gathered}$ | $\begin{gathered} -21.79 /-24.66 \\ (0.38 / 0.37) \end{gathered}$ | $\begin{gathered} -23.30 /-24.98 \\ (0.39 / 0.38) \end{gathered}$ | $\begin{gathered} 0.79 / 0.87 \\ (0.30 / 0.31) \end{gathered}$ | $\begin{gathered} 21.5 / 16.5 \\ (0.32 / 0.31) \end{gathered}$ | $\begin{gathered} 1.12 / 1.08 \\ (0.31 / 0.30) \end{gathered}$ | $\begin{aligned} & 2.24 / 2.47 \\ & 0.50 / 0.48 \end{aligned}$ | 0.12/0.09 | 10.46/10.46 | $\begin{gathered} 1.192 / 1.467 \\ (0.003 / 0.003) \end{gathered}$ | 0.30/0.28 | Y |
| RXSJ12002-0046 | J/K | $\begin{gathered} -22.51 /-24.01 \\ (0.35 / 0.37) \end{gathered}$ | $\begin{gathered} -22.79 /-24.22 \\ (0.38 / 0.36) \end{gathered}$ | $\begin{gathered} -22.99 /-24.16 \\ (0.38 / 0.38) \end{gathered}$ | $\begin{gathered} 1.39 / 1.16 \\ (0.30 / 0.31) \end{gathered}$ | $\begin{gathered} 18.5 / 16.8 \\ (0.30 / 0.30) \end{gathered}$ | $\begin{gathered} 1.10 / 1.08 \\ (0.30 / 0.30) \end{gathered}$ | $\begin{aligned} & 4.50 / 4.15 \\ & 0.30 / 0.31 \end{aligned}$ | - | - | $\begin{gathered} 1.170 / 191 \\ (0.002 / 0.002) \end{gathered}$ | 0.38/0.38 | Y |
| J12176+6546 | J/K | $\begin{gathered} -24.75 /-26.86 \\ (0.37 / 0.38) \end{gathered}$ | >-23.56/>-24.65 | (0.380. | ( | ( |  | - | - | - | $\begin{gathered} 1.241 / 1.558 \\ (0.002 / 0.002) \end{gathered}$ | 1.00/1.00 | N |
| $4 \mathrm{C}+04.42^{\text {a }}$ | J | $\begin{gathered} -27.93 \\ (0.37) \end{gathered}$ | $\begin{gathered} >-25.88 \\ - \end{gathered}$ | - | - | - | - | - | - | - | $\begin{gathered} 1.200 \\ (0.002) \end{gathered}$ | 1.00 | N |
| RXSJ $12278+3215$ | J/K | $\begin{gathered} -24.44 /-26.24 \\ (0.38 / 0.38) \end{gathered}$ | >-23.10/>-25.94 | - | - | - | ${ }^{-}$ | - | - | - | $\begin{gathered} 2.030 \\ (0.001) \end{gathered}$ | 1.00 | N |
| SDSS J124634.65+023809 ${ }^{\text {a }}$ | K | $\begin{gathered} -25.01 \\ (0.41) \end{gathered}$ | $\begin{gathered} -24.44 \\ (0.48) \end{gathered}$ | $\begin{gathered} -25.02 \\ (0.39) \end{gathered}$ | $\begin{gathered} 2.85 \\ (0.73) \end{gathered}$ | $\begin{gathered} 18.4 \\ (0.78) \end{gathered}$ | $\begin{gathered} 0.81 \\ (0.70) \end{gathered}$ | $\begin{gathered} 6.28 \\ (0.30) \end{gathered}$ | - | - | $\begin{gathered} 1.188 \\ (0.001) \end{gathered}$ | 0.36 | Y |
| J13374+6005 | J/K | $\begin{gathered} -22.84 /-24.49 \\ (0.37 / 0.39) \end{gathered}$ | $\begin{gathered} -21.88 /-23.54 \\ (0.39 / 0.41) \end{gathered}$ | $\begin{gathered} -22.30 /-23.15 \\ (0.39 / 0.40) \end{gathered}$ | $\begin{gathered} 1.11 / 1.15 \\ (0.28 / 0.30) \end{gathered}$ | $\begin{gathered} 19.5 / 17.9 \\ (0.29 / 0.31) \end{gathered}$ | $\begin{gathered} 1.10 / 1.00 \\ (0.32 / 0.35) \end{gathered}$ | $\begin{gathered} 2.42 / 3.16 \\ (0.30 / 0.34) \end{gathered}$ | 0.14/- | 11.54/- | $\begin{gathered} 1.167 / 1.205 \\ (0.001 / 0.002) \end{gathered}$ | 0.36 | N |
| J14033+0222 | J/K | $\begin{gathered} -21.91 /-24.18 \\ (0.37 / 0.38) \end{gathered}$ | $\begin{gathered} -22.98 /-24.38 \\ (0.39 / 0.39) \end{gathered}$ | $\begin{gathered} -22.89 /-24.60 \\ (0.41 / 0.42) \end{gathered}$ | $\begin{gathered} 3.98 / 3.39 \\ (0.33 / 0.35) \end{gathered}$ | $\begin{gathered} 19.7 / 18.8 \\ (0.30 / 0.30) \end{gathered}$ | $\begin{gathered} 1.11 / 1.05 \\ (0.30 / 0.30) \end{gathered}$ | $\begin{aligned} & 7.26 / 6.70 \\ & 0.37 / 0.35 \end{aligned}$ | 0.27/0.27 | 13.29/12.51 | $\begin{gathered} 1.168 / 1.172 \\ (0.001 / 0.001) \end{gathered}$ | 0.37 | N |
| $\mathrm{J} 142106+385522^{\text {a }}$ | K | $\begin{gathered} -26.63 \\ (0.35) \end{gathered}$ | $\begin{gathered} -25.68 \\ (0.37) \end{gathered}$ | $\begin{gathered} -25.94 \\ (0.41) \end{gathered}$ | $\begin{gathered} 2.89 \\ (0.46) \end{gathered}$ | $\begin{gathered} 19.3 \\ (0.42) \end{gathered}$ | $\begin{gathered} 1.11 \\ (0.38) \end{gathered}$ | $\begin{aligned} & 25.08 \\ & (0.57) \end{aligned}$ | - | - | $\begin{gathered} 1.185 \\ (0.001) \end{gathered}$ | 0.31 | Y |
| J14506+5919 | J/K | $\begin{gathered} -22.45 /-24.77 \\ (0.36 / 0.39) \end{gathered}$ | $\begin{gathered} -21.90 /-24.14 \\ (0.39 / 0.40) \end{gathered}$ | $\begin{gathered} -22.80 /-25.01 \\ (0.39 / 0.40) \end{gathered}$ | $\begin{gathered} 0.74 / 0.56 \\ (0.37 / 0.39) \end{gathered}$ | $\begin{aligned} & 20.5 / 16.00 \\ & (0.31 / 0.30) \end{aligned}$ | $\begin{gathered} 1.10 / 1.15 \\ (0.30 / 0.39) \end{gathered}$ | $\begin{gathered} 4.40 / 4.84 \\ (0.39 / 0.39) \end{gathered}$ | 0.22/0.22 | 7.98/9.14 | $\begin{gathered} 1.161 / 1.175 \\ (0.001 / 0.001) \end{gathered}$ | 0.34/0.32 | Y |
| B31441 $+476^{\text {a }}$ | K | $\begin{gathered} -28.48 \\ (0.40) \end{gathered}$ | $>-27.66$ | - | - | - | - | - | - | - | $\begin{gathered} 1.276 \\ (0.002) \end{gathered}$ | 1.00 | N |

Table 2 - continued

| Name <br> (1) | Band <br> (2) | $M_{\mathrm{nu}}$ <br> (3) | $M_{\text {bulge }}$ (4) | $\begin{gathered} M_{\text {disc }} \\ (5) \end{gathered}$ | $R \mathrm{e}(\mathrm{kpc})$ <br> (6) | $\begin{aligned} & \mu_{\mathrm{e}} \\ & (7) \end{aligned}$ | (8) | $\begin{aligned} & r_{\mathrm{s}} \\ & (9) \end{aligned}$ | $\begin{aligned} & f_{\text {bar }} \\ & (10) \end{aligned}$ | $r_{\mathrm{bar}}(\mathrm{kpc})$ (11) | $\begin{aligned} & \text { xbest } \\ & (12) \end{aligned}$ | $\chi_{\text {best }} / \chi_{\text {psf }}$ <br> (13) | $\begin{gathered} \text { Int } \\ \text { (14) } \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| PKS 1502+036 ${ }^{\text {a,c }}$ | J/K | $\begin{gathered} -24.17 /-25.12 \\ (0.38 / 0.39) \end{gathered}$ | $\begin{gathered} -23.80 /-25.96 \\ (0.39 / 0.38) \end{gathered}$ | $\begin{gathered} -24.28 /-26.30 \\ (0.39 / 0.41) \end{gathered}$ | $\begin{gathered} 1.05 / 0.49 \\ (0.64 / 0.70) \end{gathered}$ | $\begin{gathered} 15.5 / 15.7 \\ (0.36 / 0.39) \end{gathered}$ | $\begin{aligned} & 1.06 / 1.15 \\ & (0.60 / 76) \end{aligned}$ | $\begin{gathered} 4.11 / 4.17 \\ (0.30 / 0.30) \end{gathered}$ | - | - | $\begin{gathered} 1.577 / 1.236 \\ (0.001 / 0.001) \end{gathered}$ | 0.35/0.34 | Y |
| SBS 1517+520 | J/K | $\begin{gathered} -24.51 /-27.42 \\ (0.39 / 0.38) \end{gathered}$ | $\begin{gathered} -24.33 /-26.54 \\ (0.38 / 0.39) \end{gathered}$ | $\begin{gathered} -24.04 /-26.32 \\ (0.37 / 0.37) \end{gathered}$ | $\begin{gathered} 2.44 / 2.44 \\ (0.32 / 0.33) \end{gathered}$ | $\begin{gathered} 19.1 / 18.4 \\ (0.32 / 0.32) \end{gathered}$ | $\begin{gathered} 1.05 / 1.10 \\ (0.32 / 0.32) \end{gathered}$ | $\begin{gathered} 8.66 / 5.75 \\ (0.30 / 0.33) \end{gathered}$ | 0.09/0.09 | 10.26/10.26 | $\begin{gathered} 1.201 / 1.241 \\ (0.002 / 0.002) \end{gathered}$ | 0.35/0.34 | Y |
| B2 1546+35A | J/K | $\begin{gathered} -24.78 /-26.38 \\ (0.32 / 0.30) \end{gathered}$ | $\begin{gathered} -24.19 /-24.74 \\ (0.29 / 0.32) \end{gathered}$ | $\begin{gathered} -23.12 /-25.69 \\ (0.31 / 0.31) \end{gathered}$ | $\begin{gathered} 0.96 / 2.06 \\ (0.84 / 0.33) \end{gathered}$ | $\begin{gathered} 16.5 / 18.8 \\ (0.90 / 0.32) \end{gathered}$ | $\begin{gathered} 1.18 / 0.98 \\ (0.62 / 0.41) \end{gathered}$ | $\begin{aligned} & 3.99 / 10.52 \\ & (0.41 / 0.41) \end{aligned}$ | - | - | $\begin{gathered} 1.161 / 1.153 \\ (0.001 / 0.001) \end{gathered}$ | 0.32 | Y |
| J16290+4007 | J/K | $\begin{gathered} -23.85 /-25.86 \\ (0.37 / 0.39) \end{gathered}$ | >-22.64/-24.35 | - | - | - | - | - | - | - | $\begin{gathered} 1.118 \\ (0.002) \end{gathered}$ | 1.00 | Y |
| RXSJ16333+4718 | J/K | $\begin{gathered} -22.90 /-24.40 \\ (0.37 / 0.37) \end{gathered}$ | $\begin{gathered} -22.35 /-23.22 \\ (0.39 / 0.38) \end{gathered}$ | $\begin{gathered} -23.01 /-24.15 \\ (0.39 / 0.39) \end{gathered}$ | $\begin{gathered} 3.61 / 0.50 \\ (0.31 / 1.58) \end{gathered}$ | $\begin{gathered} 22.5 / 17.5 \\ (0.32 / 1.40) \end{gathered}$ | $\begin{gathered} 1.10 / 1.21 \\ (0.35 / 0.55) \end{gathered}$ | $\begin{aligned} & 8.50 / 2.83 \\ & 0.34 / 0.34 \end{aligned}$ | 0.14/0.12 | 5.04/5.04 | $\begin{gathered} 1.299 / 1.310 \\ (0.001 / 0.001) \end{gathered}$ | 0.32 | Y |
| 2E1640+5345 | J/K | $\begin{gathered} -21.69 /-23.73 \\ (0.40 / 0.39) \end{gathered}$ | $\begin{gathered} -23.40 /-24.52 \\ (0.37 / 0.37) \end{gathered}$ | $\begin{gathered} -24.21 /-25.14 \\ (0.39 / 0.39) \end{gathered}$ | $\begin{gathered} 0.45 / 0.64 \\ (0.32 / 0.32) \end{gathered}$ | $\begin{gathered} 17.5 / 16.6 \\ (0.40 / 0.38) \end{gathered}$ | $\begin{gathered} 1.10 / 1.09 \\ (0.32 / 0.32) \end{gathered}$ | $\begin{gathered} 4.50 / 5.72 \\ (0.31 / 0.31) \end{gathered}$ | 0.07/0.08 | 7.65/6.17 | $\begin{gathered} 1.178 / 1.133 \\ (0.001 / 0.001) \end{gathered}$ | 0.35 | N |
| J16410+3454 | J/K | $\begin{gathered} -21.27 /-24.35 \\ (0.39 / 0.38) \end{gathered}$ | $\begin{gathered} -23.55 /-24.65 \\ (0.39 / 0.39) \end{gathered}$ | $\begin{gathered} -24.18 /-25.55 \\ (0.38 / 0.39) \end{gathered}$ | $\begin{gathered} 0.83 / 0.98 \\ (0.49 / 0.46) \end{gathered}$ | $\begin{gathered} 17.5 / 16.8 \\ (0.34 / 0.38) \end{gathered}$ | $\begin{gathered} 2.24 / 1.53 \\ (0.43 / 0.42) \end{gathered}$ | $\begin{gathered} 4.34 / 4.01 \\ (0.35 / 0.33) \end{gathered}$ | 0.22/0.22 | 21.20/21.10 | $\begin{gathered} 1.192 / 1.169 \\ (0.001 / 0.001) \end{gathered}$ | 0.30 | Y |
| FBQSJ1644+2619 ${ }^{\text {a }}$ | J/K | $\begin{gathered} -22.51 /-24.80 \\ (0.43 / 0.24) \end{gathered}$ | $\begin{gathered} -21.41 /-24.21 \\ (0.40 / 0.32) \end{gathered}$ | $\begin{gathered} -22.90 /-24.85 \\ (0.35 / 0.38) \end{gathered}$ | $\begin{gathered} 0.96 / 1.10 \\ (0.32 / 0.34) \end{gathered}$ | $\begin{gathered} 17.5 / 16.5 \\ (0.30 / 0.31) \end{gathered}$ | $\begin{gathered} 1.80 / 1.90 \\ (1.31 / 1.35) \end{gathered}$ | $\begin{gathered} 6.65 / 7.68 \\ (0.45 / 1.14) \end{gathered}$ | -/0.17 | -/8.13 | $\begin{gathered} 1.160 \\ (0.027) \end{gathered}$ | 0.37 | Y |
| B31702+457 | K | $\begin{gathered} -23.88 \\ (0.28) \end{gathered}$ | $\begin{gathered} -25.09 \\ (0.32) \end{gathered}$ | $\begin{gathered} -23.63 \\ (0.30) \end{gathered}$ | $\begin{gathered} 1.21 \\ (0.36) \end{gathered}$ | $\begin{gathered} 17.2 \\ (0.43) \end{gathered}$ | $\begin{gathered} 1.03 \\ (0.41) \end{gathered}$ | $\begin{gathered} 5.97 \\ (0.29) \end{gathered}$ | - | - | $\begin{gathered} 2.021 \\ (0.003) \end{gathered}$ | 0.35 | N |
| J17221+5654 | J/K | $\begin{gathered} -23.84 /-25.64 \\ (0.35 / 0.36) \end{gathered}$ | $\begin{gathered} -23.95 /-24.95 \\ (0.38 / 0.39) \end{gathered}$ | $\begin{gathered} -23.12 /-24.35 \\ (0.39 / 0.39) \end{gathered}$ | $\begin{gathered} 0.71 / 1.11 \\ (0.34 / 0.36) \end{gathered}$ | $\begin{gathered} 18.3 / 17.2 \\ (0.40 / 0.35) \end{gathered}$ | $\begin{gathered} 2.10 / 2.12 \\ (0.33 / 0.32) \end{gathered}$ | $\begin{gathered} 4.34 / 4.01 \\ (0.33 / 0.32) \end{gathered}$ | - | - | $1.164 / 1.172$ $(0.001 / 0.001)$ | 0.27/0.30 | Y |
| PKS2004-447 ${ }^{\text {a,c }}$ | J/K | $\begin{gathered} -22.56 /-24.79 \\ (0.32 / 0.25) \end{gathered}$ | $\begin{gathered} -22.91 /-24.50 \\ (0.29 / 0.28) \end{gathered}$ | $\begin{gathered} -23.05 /-24.40 \\ (0.27 / 0.30) \end{gathered}$ | $\begin{gathered} 0.72 / 0.53 \\ (0.34 / 0.42) \end{gathered}$ | $\begin{gathered} 16.6 / 15.5 \\ (0.20 / 0.19) \end{gathered}$ | $\begin{gathered} 1.15 / 1.08 \\ (0.10 / 0.35) \end{gathered}$ | $\begin{gathered} 3.65 / 2.51 \\ (0.32 / 0.32) \end{gathered}$ | -/0.13 | -7.59 | $\begin{gathered} 1.403 \\ (0.025) \end{gathered}$ | 0.29 | Y |
| IRAS22453-1744 | J | $\begin{gathered} -23.34 \\ (0.35) \end{gathered}$ | $\begin{gathered} -24.10 \\ (0.36) \end{gathered}$ | - | $\begin{gathered} 2.28 \\ (0.29) \end{gathered}$ | $\begin{gathered} 17.5 \\ (0.35) \end{gathered}$ | $\begin{gathered} 0.90 \\ (0.30) \end{gathered}$ | - | - | - | $\begin{gathered} 1.176 \\ (0.001) \end{gathered}$ | 0.27 | Y |


 length; (12) the reduced $\chi^{2}$ for the best-fitting model; (13) the ratio between best-fitting $\chi^{2}$ and the PSF-fit $\chi^{2}$; (14) if the galaxy shows some sign of interaction (Y), otherwise (N).
${ }^{b}$ Galaxies with unreliable fittings according to our simulations (see Section 3.3). ${ }^{c}$ Galaxies observed using the ISAAC on the ESO/VLT.


Figure 3. The Kormendy relation in $K$ band. The symbols are explained in the figure. The bulge effective radius $\left(\log R_{\mathrm{e}}\right)$ is plotted versus the surface brightness at that radius $\left(\mu_{\mathrm{e}}\right)$. A typical error bar is shown in the lower left corner. Overlayed are a sample of high-luminosity blazars (FSRQs) and low-luminosity blazars (BL Lacs objects) from Olguín-Iglesias et al. (2016), all of them hosted by giant elliptical galaxies. For the blazars sample, we show the 95 per cent prediction bands (dotted lines) and the best linear fit (black dashed line) relation. For the NLSy1s (the 20 detected galaxies in our sample with $K$-band observations, missing are two with unreliable fittings, six undetected galaxies, and one without $K$-band observations), we show the best linear fit (yellow dashed line). There are no NLSy1s inside the 95 per cent prediction bands of the Kormendy relation, suggesting that they are hosted by disc-like bulges. A cosmological surface brightness dimming correction of the form $(1+z)^{-4}$ was applied to all targets (in this work and in Olguín-Iglesias et al. 2016).

In Table 2, we summarize the nuclear and (discy-) bulge properties. However, in Figure B1, the detailed fitting results are shown individually for every host in the sample. Out of the sample of 29 radio-loud NLSy1 host galaxies, 21 were successfully detected, seven of which are $\gamma$-rays emitters, and 14 are only radio-loud. For the sample, we estimate an average $J$-band absolute nuclear and host magnitudes of $\mathrm{M}(\mathrm{J})$ nuclear $=-23.8 \pm 1.7$ and $\mathrm{M}(\mathrm{J})$ bulge $=$ $-23.1 \pm 1.0$ and an average $K$-band absolute nuclear and host magnitudes of $\mathrm{M}(\mathrm{K})$ nuclear $=-25.6 \pm 1.4$ and $\mathrm{M}(\mathrm{K})$ bulge $=$ $-24.8 \pm 1.1$. The host luminosities for these galaxies are slightly fainter than elliptical galaxies hosting other types of radio-loud AGNs (e.g. Olguín-Iglesias et al. 2016, M(K) $)_{\mathrm{FSRQs}}=-26.2 \pm 0.9$ and $\mathrm{M}(\mathrm{K}) \mathrm{BLLacs}=-25.6 \pm 0.6)$ and rather similar to those of an $L^{*}$ galaxy $(\mathrm{M}(\mathrm{K})=-24.8 \pm 0.3$, Mobasher, Sharples \& Ellis 1993).

The average bulge effective radius for $J$ band and $K$ band, respectively, is $R_{\text {eff }}=1.9 \pm 1.3 \mathrm{kpc}$ and $R_{\text {eff }}=1.5 \pm 1.1 \mathrm{kpc}$ and the average Sérsic index is $n=1.2 \pm 0.4$ for $J$ band and $n=1.3 \pm 0.3$ for $K$ band. The Sérsic index can be used as an approximation to classify bulges (either discy, $n<2$ or classical, $n \geq$ 2, Fisher \& Drory 2008). According to the Sérsic index derived from the photometrical analysis (with the exception of $\mathrm{J} 16410+3454$ and $\mathrm{J} 17221+5654$, whose Sérsic index might be $n>2$ given their uncertainties), all the NLSy1s hosts bulges are disc-like. However, it is well known that using the Sérsic index to discern between classical from disc-like bulges can generate many misclassifications (Gadotti 2009). To address this, we used the fact that disc-like and classical bulges are expected to be structurally different. Therefore, disc-like bulges should not follow the Kormendy relation (inverse relation between $\mu_{\mathrm{e}}$ and $R_{\mathrm{e}}$ found in elliptical galaxies and classical bulges; Kormendy 1977). In fact, Gadotti (2009) finds that disc-like bulges lie below the Kormendy relation and hence it can be used to distinguish them from classical bulges.

In Fig. 3, we explore the Kormendy relation by plotting the effective radius $\left(R_{\mathrm{e}}\right)$ versus the surface brightness at the bulge ef-
fective radius ( $\mu_{\mathrm{e}}$ ) of the detected hosts with $K$-band observations. ${ }^{3}$ Together, we plot the results of the blazars hosts analysis from Olguín-Iglesias et al. (2016). The blazars hosts show a statistically significant correlation between $R_{\mathrm{e}}$ and $\mu_{\mathrm{e}}$ (the Kormendy relation). On the other hand, the NLSy1s in the sample show a shallower trend. The 95 per cent prediction bands of the correlation are represented by the dotted lines. Since the aim of the prediction bands is to encompass the likely values of future observations from the same sampled population, we might say that it is most likely that the hosts that lie below the lower prediction band are not classical bulges. We can see that all ( 20 with host detections, $K$-band observations and inside the 'correctly retrieved' area) RL-NLSy1s lie below the Kormendy relation and below the 95 per cent prediction bands. If errors are taken into account, four NLSy1s might be consistent with classical bulges and 16 are certainly disc-like bulges.

Although these results suggest disc-like systems as RL-NLSy1s hosts, some RL-NLSy1s might be hosted by classical-bulges. Therefore, further studies of their stellar populations and kinematics, using integral field spectroscopy, will help in understanding their nature. This result is not surprising since the prevalence of disclike bulges in this type of AGN hosts have been previously found (e.g. Orban de Xivry et al. 2011; Mathur et al. 2012). However, very little is known about the presence of radio jets launched from disc galaxies. In the past, only a small number of disc-like systems were found to be radio galaxies (e.g. Ledlow, Owen \& Keel 1998; Hota et al. 2011; Morganti et al. 2011; Bagchi et al. 2014; Kaviraj et al. 2015; Mao et al. 2015; Singh et al. 2015). However, only three (included in this work) have been found to, additionally, be $\gamma$-rays emitters (León Tavares et al. 2014; Kotilainen et al. 2016; Olguín-Iglesias et al. 2017, 1H0323 + 342, PKS2004-447, and FBQSJ1644+2619, respectively).

The uniqueness of the hosts of NLSy1s (and nearby environment) might hint at how these galaxies acquired such properties. Spiral arms, bar incidence, interactions evidence, etc., could shed some light on the fuelling mechanisms needed by the central SMBH to form and develop powerful relativistic jets. Therefore, in the following sections we discuss on the specific features that the hosts of the galaxies in the sample show.

### 4.1 Bar frequency

Through a simple visual inspection, some galaxies reveal the presence of a bar in their brightness distributions. However, in order to have a quantitative identification of these bars, we adopt an analysis based on the ellipse fit to the galaxy isophotes, where radial variations of ellipticity ( $\epsilon$ ) and position angle (PA) exhibit the existence, ellipticity, and extent of the bar. Moreover, some bars detected by this way might not be obvious at a glance. The detection of bars using this method consists of finding a local maximum in $\epsilon$, which indicates the bar end. Along the bar, the PA should remain constant, thus along the suspected bar the PA should not change much (typically $\Delta \mathrm{PA} \lesssim 20^{\circ}$, Wozniak et al. 1995; Jogee, Kenney \& Smith 1999; Menéndez-Delmestre et al. 2007). At larger radius, further away from the bar end, we should measure the ellipticity and PA of the disc, then the ellipticity should drop (at least 0.1 ; $\Delta \epsilon \geq 0.1$ ) and most likely the PA will change. In Figure B2, we show the plots of PA and ellipticity versus radius of 10 of the galaxies that fulfil this criteria (the other two galaxies in the sample

[^3]with detected bars are shown in Kotilainen et al. 2016 and OlguínIglesias et al. 2017). Thus, the detected hosts with bars represent 52 per cent ( 77 per cent RL-NLSy1s and 20 per cent $\gamma$-NLSy1s) of the galaxies in the sample, whereas Laurikainen, Salo \& Buta (2004) find bars in 62 per cent $\pm 9$ per cent of Seyfert galaxies and Järvelä, Lähteenmäki \& Berton (2018) find bars in four out of five sources of a sample of NLSy1s that could launch powerful relativistic jets. It must be stressed that, although near-IR imaging (specially $K$ band) provides a reliable assessment of the bar fraction, we might miss some bars given the relative high redshift of some of the galaxies in the sample (e.g. J17221+5654 is the galaxy with the highest redshift $z=0.426$ and a bar detection).

In addition to finding bars, the radial variations of ellipticity and PA help in estimating their ellipticity and length. Ellipticity has been shown to be a good bar strength indicator (Laurikainen, Salo \& Rautiainen 2002). Abraham \& Merrifield (2000) defined a bar strength parameter given by
$\mathrm{f}_{\text {bar }}=\frac{2}{\pi}\left(\arctan \left(1-\epsilon_{\text {bar }}\right)^{-1 / 2}-\arctan \left(1-\epsilon_{\text {bar }}\right)^{1 / 2}\right)$,
where $\epsilon_{\text {bar }}$ is the bar ellipticity at the bar end. Using this parameter, we find that, for our sample, $\mathrm{f}_{\mathrm{bar}}=0.13 \pm 0.06$. By comparing this value with the findings of Laurikainen et al. (2007; average $f_{\text {bar }}=$ $0.20 \pm 0.03$, from 216 galaxies observed in the NIR, including all Hubble types), we note that the bar strengths for our NLSy1s sample is rather low. By contrast, we find that the average bar length of our sample ( $r_{\text {bar }}=9.8 \mathrm{kpc} \pm 1.8$ ) is large when compared either with late-type or early-type galaxies ( $\sim 0.5-3.5 \mathrm{kpc}$ and $\sim 1-10 \mathrm{kpc}$, respectively; Erwin 2005). We note that, even if we assume that the galaxies with no bars detected in our sample have the shortest bars ( 0.5 kpc ), the average bar length would be 5.9 kpc , still larger than the average for late-type galaxies ( 2 kpc ), early-type galaxies ( 5.5 kpc ), or both early- and late-type galaxies ( 3.75 kpc ). These results might hint to the necessary conditions to properly channel the fuel towards the centre of the galaxy to feed the black hole and trigger its activity. According to bar models by Athanassoula (1992), weak bars promote the inflow of gas towards the inner Lindblad resonance (ILR), and forms a nuclear ring. As long as the bar pattern speed remains low, the ILR is kept, which occurs provided that bars are long. If dynamical instabilities via gradual build-up of material show, material from that nuclear ring would flow inwards and trigger the black hole activity, as hypothesized by Laurikainen et al. (2002). Low and steady evolution (secular evolution) might thus be capable of producing AGNs as powerful as to launch radio jets.

### 4.2 Mergers and galaxy interactions

Secular evolution driven by stellar asymmetries (i.e. bars, lopsidedness, spiral patterns, and other coherent structures) can be strengthened by external processes such as tidal interactions and mergers (Mapelli, Moore \& Bland-Hawthorn 2008; Reichard et al. 2009). In our sample, 62 per cent of the galaxies show some sign of merger, interaction, or off-centred components (nine $\gamma$-NLSy1s and nine RL-NLSy1s, including both detected and non-detected hosts). This result is important since both observations and simulations suggest that AGN activity is closely related to galaxy interactions and mergers (e.g. Di Matteo et al. 2003; Ellison et al. 2011; Silverman et al. 2011; Koss et al. 2012; Ellison et al. 2013; Capelo et al. 2015; Järvelä et al. 2018).

Particularly, in the case of the $\gamma$-NLSy1s, we note that only three out of $12 \gamma$-NLSy1s do not show signs of interac-

Table 3. Summary of galaxies with offset AGNs.

| Galaxy | Redshift | $\Delta x(\operatorname{arcsec})$ | $\Delta x(\mathrm{kpc})$ |
| :--- | :---: | :---: | :---: |
| RXSJ12002-0046 | 0.179 | 0.24 | 0.73 |
| J142106+385522 | 0.490 | 0.25 | 1.50 |
| J14506+5919 | 0.202 | 0.23 | 0.76 |
| B2 1546+35A | 0.479 | 0.24 | 1.42 |
| J17221+5654 | 0.426 | 0.23 | 1.30 |
| IRAS 22453-1744 | 0.117 | 0.71 | 1.50 |

Notes. Column (1) gives the galaxy name; (2) the redshift of the system; (3) and (4) the projected separation between the stellar bulge and the AGN in arcsec and kpc, respectively. The typical error is $\pm 0.16$ arcsec, derived as described in Section 3.2.2.
${ }^{a} \gamma$-ray emitting NLSy1 galaxy.
tions (PMNJ0948+0022, $z=0585 ; 4 \mathrm{C}+04.42, z=0.996 ; \mathrm{B} 3$ $1441+476, z=705$ ), considering their high redshift and that only one of these hosts is detected. The large fraction of interactions in the $\gamma$-NLSy1s of the sample ( 75 per cent against 53 per cent, when compared with RL-NLSy1s) largely consists of the host itself and another, significantly smaller, galaxy or faint tidal feature (i.e. minor mergers; however, spectroscopic data are required to confirm the idea of these features as interactions).

This result is important since simulations (e.g. Qu et al. 2011) show that the angular momentum decreases more significantly when the stellar disc undergoes a minor merger than when it evolves in isolation. Hence, the difference in power between RL-NLSy1s and $\gamma$-NLSy1s might thus be the result of the difference between the processes that drive their evolution. In this way, our findings suggest secular evolution as a process capable of producing not only radio, but also $\gamma$-ray emitting jets.

Another important finding in our study is that not only dises might be off-centred with respect to the nucleus, also some bulges might. This behaviour has not only been predicted by simulations (Hopkins et al. 2012) but also previously observed. The first offset AGN reported was the Seyfert 1 galaxy NGC 3227 (Mediavilla \& Arribas 1993), where the region of broad emission lines is offset from the kinematic centre by $\sim 0.250 \mathrm{kpc}$. Another important example is the low-luminosity AGN NGC 3115 (Menezes, Steiner \& Ricci 2014), with an AGN located at a projected distance of $\sim 0.014 \mathrm{kpc}$ from the stellar bulge. Similarly, using GALFIT and observations from ChandralACIS and the Hubble Space Telescope, Comerford et al. (2015) analysed a sample of 12 dual AGN candidates at $z<0.34$ and discovered six systems that are either dual or offset AGNs with separations $\Delta x<10 \mathrm{kpc}$. Finally, here we find a total of six systems (see Table 3) where the stellar bulge is offset from the AGN by projected distances $\Delta x<1.5 \mathrm{kpc}$

This finding strongly suggests an important connection between AGNs and galaxy mergers. Two likely scenarios where the AGN is off-centred with respect to the stellar bulge are explained on the basis of galaxy mergers. On the one hand, the black hole of the observed AGN and another (inactive) black hole form a binary system. The inactive black hole is located in the centre of stellar bulge, and thus the AGN is offset with respect to it (Menezes et al. 2014). On the other hand, two black holes might have already coalesced which caused the formation of gravitational waves, which in turn asymmetrically pushed the system to shape it to its current form (Merritt 2006; Blecha \& Loeb 2008; Sundararajan, Khanna \& Hughes 2010; Blecha et al. 2019, and references therein). This important finding might help in constraining SMBH-galaxy co-evolution theoretical studies and simulations where most of the times a stationary central black hole is assumed.

Table 4. Average $J-K$ colours for the bulge, disc, and nucleus of the host galaxies in the sample.

| Subsample | $(J-K)_{\text {bulge }}$ | $(J-K)_{\text {disc }}$ | $(J-K)_{\text {nuclear }}$ |
| :--- | :---: | :---: | :---: |
| All | $1.7 \pm 0.7$ | $1.7 \pm 0.5$ | $2.0 \pm 0.5$ |
| RL-NLSy1s | $1.5 \pm 0.5$ | $1.6 \pm 0.6$ | $2.0 \pm 0.5$ |
| $\gamma$-NLSy1s | $2.1 \pm 0.5$ | $1.8 \pm 0.4$ | $1.9 \pm 0.5$ |

### 4.2.1 Notes on individual galaxies

Here, we provide a short description of the characteristics that each interacting galaxy shows (for images, see Appendix B1).
(i) $1 \mathrm{H} 0323+342$. The closest $\gamma$-ray emitter NLSy1 galaxy. In this galaxy a ringed structure is seen which is interpreted as 'suggestive evidence for a recent violent dynamical interaction' by León Tavares et al. (2014), where an extensive discussion on this galaxy can be found.
(ii) SBS $0846+513$. The host galaxy of this $\gamma$-NLSy1 is not detected, therefore it was modelled using a PSF. A bright and close companion ( $<5 \mathrm{arcsec}$ ) is clearly detected. Also, a spiral galaxy in the foreground is observed.
(iii) J093241+53063. A source with a disc-like bulge, with Sérsic index $n=1.21 \pm 0.70$. Although no disc is detected, a faint companion is.
(iv) J095820 +322401 . Although it has relatively high redshift ( $z=0.53$ ), we detect a faint close neighbour in this galaxy. When analysed further, an even fainter disc is revealed in our analysis. The disc is off-centred with respect to the nucleus and bulge. This maybe due to the action exerted by its alleged neighbour.
(v) FBQSJ1102+2239. This galaxy represents a classical encounter between two disc galaxies. Very similar both in $J$ and $K$ bands, a remnant of the disc is detected in the AGN host. The companion keeps a spiral arm and it is barely connected with the main AGN host. Another blob is observed, an H II region, which is part of the system (according to optical spectra).
(vi) IRAS 11598-0112. The AGN host is modelled using a disc-like bulge and a disc. However, two spiral-arm-like features are included in the model. More interestingly, another component (probably a discy bulge) is detected inside the main disc.
(vii) SDSS J124634.65+023808. An AGN-dominated galaxy modelled using a bulge and an exponential disc. It shows a close and faint feature which is fitted using an exponential disc.
(viii) J14033+0222. An apparent simple barred galaxy. However, the disc is off-centred with respect to the AGN and bulge (in both $J$ and $K$ bands), which suggests a non-obvious interaction. Since both bar and disc have exponential profiles and the bar is faint, only one exponential was needed to model both.
(ix) J142106+385522. The host galaxy of this $\gamma$-NLSy1 was modelled using a bulge and a disc. In the image, a faint tail-like feature (which was not modelled due to its intensity) is observed. Suggestion of an interaction are observed in its off-centred disc.
(x) PKS1502+036. A clear companion only visible in $J$ band. The companion seems close, however, this galaxy does not show asymmetries as others in the sample. Although D'Ammando et al. (2018) find it hosted in an elliptical galaxy, we find a better fit using a disc-like host.
(xi) SBS1517+520. In this galaxy, the asymmetry of its surface brightness profile and its off-centred disc (more evident in the $K$ band) hint towards some type of disruption induced by an interaction.
(xii) B2 $1546+35 \mathrm{~A}$. When the host is decomposed into bulge and disc, the different parts are off-centred with respect to each other. While the host galaxy image, both in $J$ and $K$ bands, looks similar, when it is represented using its surface brightness profile, the two bands differ.
(xiii) RXSJ16333+4718. Two disc galaxies interacting. Seemingly, the companion is also face-on. Again, the host galaxy shows a greater effect of the interaction on one of the bands. While in $K$ band, the bulge seems a bit off-centred, in the $J$ band both the disc and the bulge shows greater impact on its morphology.
(xiv) J16410+3454. This galaxy shows a feature that emerges after the fitting of a bulge + disc + AGN model. This additional component is modelled using a Sérsic profile with $\mathrm{n} \approx 1$. The main disc seems off-centred in the radial profiles (maybe because of the effect of the interaction).
(xv) FBQS J1644+2619. In $J$ band, a ring feature shows, whose formation process might be that described by Athanassoula, Puerari \& Bosma (1997) and that PKS 2004-447 might be undergoing. Besides, a faint disruption of the ring suggests an interaction. On the other hand, in $K$ band, a bar is observed and the ring features is almost absent. An extensive discussion on these features is found in Olguín-Iglesias et al. (2016).
(xvi) J17221+5654. The host galaxy was modelled using a bulge and a disc. In both $J$ and $K$ bands a small companion is detected a few arcseconds away from the host centre. The companion seems to be interacting with the main galaxy since the components of the AGN host are off-centred; more evident in $K$ band.
(xvii) PKS 2004-447. This barred galaxy shows two faint spiral arms, one of which is more open. It is a very good example of part of the evolution of the simulations by Athanassoula et al. (1997), where the impact of a small companion on a barred galaxy leads to the formation of a ring. For a detailed discussion on this galaxy see Kotilainen et al. (2016).
(xviii) IRAS 22453-1744. A bulge with a Sérsic index $n=$ $0.90 \pm 0.30$ was used to model the host. However, the AGN and host are off-centred. The most likely reason is the close (merged) companion with a complex morphology that is difficult to characterize.

### 4.3 AGN, bulge, and disc $(J-K)$ colours

Table 4 shows the average $J-K$ colours of the main components of the host galaxies in the sample (whenever they have both $J$ and $K$-band information). We see that the disc and nuclear $J-K$ colours remain virtually unchanged whether the subsample includes $\gamma$-NLSy1s or only RL-NLSy1s. We also note that, unexpectedly, bulge and disc colours for RL-NLSy1s are the same within errors (i.e. bulges are as blue as discs, when disc are expected to be bluer, Moriondo, Giovanardi \& Hunt 1998; Seigar \& James 1998). This result is thus consistent with star-forming bulges, which imply large gas reservoirs.
On the other hand, the average $J-K$ bulge colour changes depending on the subsample, being redder for $\gamma$-NLSy1s ( $J-K=$ $2.1 \pm 0.5$ ). According to findings by Glass \& Moorwood (1985) and Seigar \& James (1998), NIR colours $J-K \approx 2.0$ could be the result of a dust-embedded AGN or a nuclear starbursts (if the component is extended, i.e. in bulges). Bulge reddening might be linked to the large fraction of $\gamma$-NLSy1s showing signs of minor mergers. Thus, according to these results, interactions are likely to play an important role in the nuclear activity of the galaxies in our sample.


Figure 4. Plot of the nuclear $K$-band magnitude versus the bulge $K$-band magnitude for the host galaxies of the sample. The symbols are explained in the figure. The upper limits for unresolved galaxies are shown as down arrows. We show the 25 detected and undetected galaxies in our sample with $K$-band observations. Missing are two with unreliable fittings and two without $K$-band observations. Overlayed is a sample of blazars from Olguín-Iglesias et al. (2016), where the galaxies were analysed in $K$ band. For the blazars sample we show the best linear fits (dashed blue line for FSRQs and dashed red line for BL Lacs) and the 95 per cent prediction bands (dotted black lines). The 99 per cent confidence intervals are shown for FSRQs (blue shade) and NLSy1s (yellow shade). For the NLSy1s sample we show the best linear fit (green solid line). All $K$-corrections were performed using the $K$-corrections calculator (Chilingarian et al. 2010; Chilingarian \& Zolotukhin 2012).

## 5 AGN-HOST CONNECTION

In Fig. 4 , we explore the $M_{\mathrm{K}}$, nuclear $-M_{\mathrm{K} \text {, bulge }}$ plot for our sample. The first thing we notice is that most of the sample data points fall in the bottom left corner, where high-luminosity blazars (HLBs), i.e. flat spectrum radio quasars (FSRQs) and low-luminosity blazars (LLBs), i.e. BL Lacs coincide. However, three $\gamma$-NLSy1s (only one with host galaxy detection) are brighter and lie $\sim 3$ mag apart from the group and two (also $\gamma$-NLSy1s) are fainter and lie $\sim 1 \mathrm{mag}$ appart. At first glance the bulk of RL-NLSy1s are part of either HLB or LLB. A further analysis shows that, if we include the NLSy1s in the HLB sample, the best linear fit marginally lies inside the 99 per cent confidence intervals of the FSRQs best linear fit. There is a statistically significant positive correlation for the FSRQs sample ( $r=0.8, \mathrm{p}=4 \times 10^{-12}$ ) which is kept virtually unchanged ( $\mathrm{p} \approx 10^{-12}$ ) when the RL-NLSy1s are added to the FSRQs.

When we analyse the NLSy1s sample alone, a statistically significant positive correlation is observed ( $r=0.6, \mathrm{p}=1 \times 10^{-3}$ ). Whether the NLSy1s sample belongs to the FSRQs sample or not, its positive trend suggests that their jets also stimulate starburst activity in their hosts.

The results presented above suggest either FSRQs and RLNLSy1s accretion modes locate them close to each other in the $M_{\mathrm{K}, \text { nuclear }}-M_{\mathrm{K} \text {, bulge }}$ plot. Bearing in mind that both FSRQs and NLSy1s are thought to accrete matter very efficiently via accretion discs, this is expected.

In order to rule out the scenario where RL-NLSy1s behave as LLB/BL Lacs (since, independently of nuclear magnitude, BL Lacs show a narrow range of bulge magnitudes, where the bulk of RL-NLSy1s bulge magnitudes also reside), we perform a twodimensional Kolmogorov-Smirnov test for the bulge and nuclear
magnitudes of BL Lacs and RL-NLSy1s. According to the results of this test $\left(p \approx 3.0^{-4}\right)$ there is a statistically significant difference between the samples, suggesting that, in the $M_{\mathrm{K}, \text { nuclear }}-M_{\mathrm{K}}$, bulge plot, BL Lacs and RL-NLSy1s are not similar.

Previous studies conducted on RL-NLSy1s had already supported the idea that, when compared with blazars, RL-NLSy1s are particularly similar to FSRQs (e.g. Paliya et al. 2013; Foschini et al. 2015; Paliya \& Stalin 2016). Provided that RL-NLSy1s are hosted by spiral galaxies, what our findings suggest might be a substantive contribution since blazars are known to be hosted by elliptical galaxies and powered by massive black holes.

The positive trend in the NLSy1s and FSRQs samples clearly suggest a positive feedback scenario. The AGN outflows can induce star formation, both in the galactic disc through compression of molecular clouds (Silk 2013) or directly in the outflowing gas (Ishibashi \& Fabian 2012).

## 6 SUMMARY

We have presented near-infrared images of a sample of 29 radio-loud NLSy1 host galaxies, 12 of which are also classified as $\gamma$-NLSy1s. By thoroughly analysing their 2D surface brightness distribution, we successfully detected 21 hosts (14 RL-NLSy1s and $7 \gamma$-NLSy1s). Our near-infrared study allowed us to compare the photometrical properties of RL-NLSy1s with another type of AGN capable of launching powerful radio jets, namely blazars (both BL Lacs and FSRQs). The main findings of our study are summarized below.
(i) The photometrical properties derived by our 2D analysis for a sample of RL-NSLy1s suggests that, consistent with radio-quiet NLSy1s and opposite to the jet paradigm, powerful relativistic jets
can be launched from disc-like systems instead of elliptical galaxies or classical bulges.
(ii) Secular evolution driven by the peculiar bar properties in our sample (long and weak) might be responsible for channelling fuel towards the centre of the galaxy to feed the black hole and trigger the nuclear activity, whereas the nuclear activity in $\gamma$-NLSy 1 s could be the result of a similar process enhanced by minor mergers.
(iii) RL-NLSy1s bulges and discs show the same average NIR colour $(J-K)_{\text {disc }}=1.6 \pm 0.6$ and $(J-K)_{\text {bulge }}=1.5 \pm 0.5$. Since discs are expected to be bluer, this result is consisting with star-forming bulges, suggesting large gas reservoirs. On the other hand, $\gamma$-NLSy1s bulges show an average NIR colour $(J-K)_{\mathrm{AGN}}=$ $2.1 \pm 0.5$. This reddening (with respect to RL-NLSy1s) suggests nuclear starburts, probably linked to the large fraction of minor mergers shown by $\gamma$-NLSy 1 s , which in turn could make a difference between RL- and $\gamma$-NLSy1s nuclear activity.
(iv) We have discovered six systems showing an offset stellar bulge with respect to the AGNs (with separations $\Delta<1.5 \mathrm{kpc}$ ). This might be the result of a galaxy merger, strongly suggesting an important connection between AGNs and galaxy mergers.
(v) Hints of positive feedback are suggested when we plot $M_{\mathrm{K}, \text { nuclear }}$ versus $M_{\mathrm{K}, \text { bulge }}$ of the sample. We find that RL-NLSy1s behave in a similar manner as FSRQs (or HLBs), which might be the result of a similar accretion mode between RL-NLSy1s and FSRQs.

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## REFERENCES

Abdo A. A. et al., 2009a, ApJ, 699, 976
Abdo A. A. et al., 2009b, ApJ, 707, L142
Abraham R. G., Merrifield M. R., 2000, AJ, 120, 2835
Athanassoula E., 1992, MNRAS, 259, 345
Athanassoula E., Puerari I., Bosma A., 1997, MNRAS, 286, 284
Bagchi J. et al., 2014, ApJ, 788, 174
Bertin E., Arnouts S., 1996, A\&ASS, 117, 393
Berton M. et al., 2019, AJ, 157, 48
Blecha L., Loeb A., 2008, MNRAS, 390, 1311
Blecha L., Brisken W., Burke-Spolaor S., Civano F., Comerford J., Darling J., Lazio T. J. W., Maccarone T. J., 2019, Astro2020: Decadal Survey on Astronomy and Astrophysics, 2020, 318
Capelo P. R., Volonteri M., Dotti M., Bellovary J. M., Mayer L., Governato F., 2015, MNRAS, 447, 2123

Chilingarian I. V., Zolotukhin I. Y., 2012, MNRAS, 419, 1727
Chilingarian I. V., Melchior A.-L., Zolotukhin I. Y., 2010, MNRAS, 405, 1409
Comerford J. M., Pooley D., Barrows R. S., Greene J. E., Zakamska N. L., Madejski G. M., Cooper M. C., 2015, ApJ, 806, 219
D'Ammando F. et al., 2012, MNRAS, 426, 317
D’Ammando F., Orienti M., Larsson J., Giroletti M., 2015, MNRAS, 452, 520
D’Ammando F., Acosta-Pulido J. A., Capetti A., Baldi R. D., Orienti M., Raiteri C. M., Ramos Almeida C., 2018, MNRAS, 478, L66
Di Matteo T., Croft R. A. C., Springel V., Hernquist L., 2003, ApJ, 593, 56
Ellison S. L., Patton D. R., Mendel J. T., Scudder J. M., 2011, MNRAS, 418, 2043
Ellison S. L., Mendel J. T., Patton D. R., Scudder J. M., 2013, MNRAS, 435, 3627

Erwin P., 2005, MNRAS, 364, 283
Fabian A. C., 2012, ARA\&A, 50, 455
Ferrarese L., Merritt D., 2000, ApJ, 539, L9
Fisher D. B., Drory N., 2008, in Funes J. G., Corsini E. M., eds, ASP Conf. Ser. Vol. 396, Formation and Evolution of Galaxy Discs. Astron. Soc. Pac., San Francisco, p. 309
Foschini L., 2011, Proc. Sci., Narrow-Line Seyfert 1 Galaxies and their Place in the Universe, SISSA, Trieste, PoS\#24
Foschini L. et al., 2015, A\&A, 575, A13
Gadotti D. A., 2009, MNRAS, 393, 1531
Gebhardt K. et al., 2000, ApJ, 539, L13
Glass I. S., Moorwood A. F. M., 1985, MNRAS, 214, 429
Gültekin K. et al., 2009, ApJ, 698, 198
Heckman T. M., Best P. N., 2014, ARA\&A, 52, 589
Hopkins P. F., Hernquist L., Hayward C. C., Narayanan D., 2012, MNRAS, 425, 1121
Hota A. et al., 2011, MNRAS, 417, L36
Hyvönen T., Kotilainen J. K., Falomo R., Örndahl E., Pursimo T., 2007, A\&A, 476, 723
Ishibashi W., Fabian A. C., 2012, MNRAS, 427, 2998
Järvelä E., Lähteenmäki A., Berton M., 2018, A\&A, 619, A69
Jogee S., Kenney J. D. P., Smith B. J., 1999, ApJ, 526, 665
Kaspi S., Smith P. S., Netzer H., Maoz D., Jannuzi B. T., Giveon U., 2000, ApJ, 533, 631
Kaviraj S., Shabala S. S., Deller A. T., Middelberg E., 2015, MNRAS, 454, 1595
Komossa S., Voges W., Xu D., Mathur S., Adorf H.-M., Lemson G., Duschl W. J., Grupe D., 2006, AJ, 132, 531

Kormendy J., 1977, ApJ, 218, 333
Koss M., Mushotzky R., Treister E., Veilleux S., Vasudevan R., Trippe M., 2012, ApJ, 746, L22
Kotilainen J. K., Falomo R., Scarpa R., 1998a, A\&A, 332, 503
Kotilainen J. K., Falomo R., Scarpa R., 1998b, A\&A, 336, 479
Kotilainen J. K., Hyvönen T., Falomo R., 2005, A\&A, 440, 831
Kotilainen J. K., Falomo R., Labita M., Treves A., Uslenghi M., 2007, ApJ, 660, 1039
Kotilainen J. K., Tavares J. L., Olguin-Iglesias A., Baes M., Anorve C., Chavushyan V., Carrasco L., 2016, ApJ, 832, 157
Laurikainen E., Salo H., Rautiainen P., 2002, MNRAS, 331, 880
Laurikainen E., Salo H., Buta R., 2004, ApJ, 607, 103
Laurikainen E., Salo H., Buta R., Knapen J. H., 2007, MNRAS, 381, 401
Ledlow M. J., Owen F. N., Keel W. C., 1998, ApJ, 495, 227
León Tavares J. et al., 2014, ApJ, 795, 58
Liao N.-H., Xin Y.-L., Li S., Jiang W., Liang Y.-F., Li X., Zhang P.-F., Chen L., et al., 2015, ApJ, 808, 74

Magorrian J. et al., 1998, AJ, 115, 2285
Mao M. Y. et al., 2015, MNRAS, 446, 4176
Mapelli M., Moore B., Bland-Hawthorn J., 2008, MNRAS, 388, 697
Mathur S., Fields D., Peterson B. M., Grupe D., 2012, ApJ, 754, 146
Mediavilla E., Arribas S., 1993, Nature, 365, 420
Menéndez-Delmestre K., Sheth K., Schinnerer E., Jarrett T. H., Scoville N. Z., 2007, ApJ, 657, 790

Menezes R. B., Steiner J. E., Ricci T. V., 2014, ApJ, 796, L13
Merritt D., 2006, ApJ, 648, 976
Mobasher B., Sharples R. M., Ellis R. S., 1993, MNRAS, 263, 560
Moorwood A. et al., 1998, The Messenger, 94, 7
Morganti R., Holt J., Tadhunter C., Ramos Almeida C., Dicken D., Inskip K., Oosterloo T., Tzioumis T., 2011, A\&A, 535, A97

Moriondo G., Giovanardi C., Hunt L. K., 1998, A\&AS, 130, 81
Ojha V., Krishna G., Chand H., 2019, MNRAS, 483, 3036
Olguín-Iglesias A. et al., 2016, MNRAS, 460, 3202
Olguín-Iglesias A., Kotilainen J. K., León Tavares J., Chavushyan V., Añorve C., 2017, MNRAS, 467, 3712

Orban de Xivry G., Davies R., Schartmann M., Komossa S., Marconi A., Hicks E., Engel H., Tacconi L., 2011, MNRAS, 417, 2721
Osterbrock D. E., Pogge R. W., 1985, ApJ, 297, 166
Paliya V. S., Stalin C. S., 2016, ApJ, 820, 52

Paliya V. S., Stalin C. S., Kumar B., Kumar B., Bhatt V. K., Pandey S. B., Yadav R. K. S., 2013, MNRAS, 428, 2450
Paliya V. S., Ajello M., Rakshit S., Mandal A. K., Stalin C. S., Kaur A., Hartmann D., 2018, ApJ, 853, L2
Peng C. Y., Ho L. C., Impey C. D., Rix H.-W., 2011, Astrophysics Source Code Library, record ascl:1104.010
Qu Y., Di Matteo P., Lehnert M. D., van Driel W., Jog C. J., 2011, A\&A, 535, A5
Reichard T. A., Heckman T. M., Rudnick G., Brinchmann J., Kauffmann G., Wild V., 2009, ApJ, 691, 1005

Scarpa R., Urry C. M., Falomo R., Pesce J. E., Treves A., 2000, ApJ, 532, 740
Seigar M. S., James P. A., 1998, MNRAS, 299, 672
Silk J., 2013, ApJ, 772, 112
Silverman J. D. et al., 2011, ApJ, 743, 2
Singh V., Ishwara-Chandra C. H., Sievers J., Wadadekar Y., Hilton M., Beelen A., 2015, MNRAS, 454, 1556
Skrutskie M. F. et al., 2006, AJ, 131, 1163
Stickel M., Padovani P., Urry C. M., Fried J. W., Kuehr H., 1991, ApJ, 374, 431
Sundararajan P. A., Khanna G., Hughes S. A., 2010, Phys. Rev. D, 81, 104009
Tremaine S. et al., 2002, ApJ, 574, 740

Urry C. M., Scarpa R., O’Dowd M., Falomo R., Pesce J. E., Treves A., 2000, ApJ, 532, 816
Wozniak H., Friedli D., Martinet L., Martin P., Bratschi P., 1995, A\&AS, 111, 1153.3
Yao S., Yuan W., Zhou H., Komossa S., Zhang J., Qiao E., Liu B., 2015, MNRAS, 454, L16

## SUPPORTING INFORMATION

Supplementary data are available at $M N R A S$ online.
Fig. SB1. Surface brightness profile decomposition of the host galaxies in the sample.
Fig. SB2. Radial variation of ellipticity $\epsilon$ (blue circles) and of PA (red squares) derived by ellipse fitting to the galaxy isophotes.

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## APPENDIX A: SIMULATIONS

Table A1. Parameters of the sample of simulated galaxies (columns 2-6) and the retrieved parameters of the simulated galaxies when they are modelled in an identical way to the real galaxies in our sample (columns 7-12).

|  |  | Simu | ation par | meters |  |  |  | Retrieved | aramet |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Seeing <br> (arcsec) | $m_{\text {nuclear }}$ | $m_{\text {bulge }}$ | $\begin{gathered} R \mathrm{e} \\ (\operatorname{arcsec}) \end{gathered}$ | $m_{\text {disc }}$ | $m_{\text {nuclear }}$ | $m_{\text {bulge }}$ | $R_{\mathrm{e}}$ <br> (arcsec) | $n$ | $m_{\text {disc }}$ | $\begin{gathered} R_{\mathrm{S}} \\ (\operatorname{arcsec}) \end{gathered}$ | Model quality |
| 1 | 0.60 | 16.00 | 18.50 | 0.10 | 15.00 | 15.90 | 20.75 | 2.13 | 0.06 | 15.09 | 2.83 | bad |
| 2 | 0.60 | 15.00 | 17.00 | 0.10 | 15.00 | 14.85 | 19.63 | 2.09 | 0.19 | 15.09 | 2.87 | bad |
| 3 | 0.60 | 14.00 | 15.50 | 0.10 | 15.00 | 13.79 | 17.71 | 0.40 | 0.06 | 15.09 | 2.79 | bad |
| 4 | 0.60 | 13.00 | 14.00 | 0.10 | 15.00 | 12.66 | 17.92 | 0.21 | 1.84 | 15.09 | 2.75 | bad |
| 5 | 0.60 | 12.00 | 12.50 | 0.10 | 15.00 | 11.64 | 13.54 | 0.25 | 0.05 | 15.09 | 2.81 | bad |
| 6 | 0.60 | 16.00 | 18.50 | 0.10 | 15.00 | 15.90 | 20.75 | 2.13 | 0.06 | 15.09 | 2.83 | bad |
| 7 | 0.60 | 15.00 | 17.00 | 0.10 | 15.00 | 14.85 | 19.63 | 2.09 | 0.19 | 15.09 | 2.87 | bad |
| 8 | 0.60 | 14.00 | 15.50 | 0.10 | 15.00 | 13.79 | 17.71 | 0.40 | 0.06 | 15.09 | 2.79 | bad |
| 9 | 0.60 | 13.00 | 14.00 | 0.10 | 15.00 | 12.66 | 17.92 | 0.21 | 1.84 | 15.09 | 2.75 | bad |
| 10 | 0.60 | 12.00 | 12.50 | 0.10 | 15.00 | 11.64 | 13.54 | 0.25 | 0.05 | 15.09 | 2.81 | bad |
| 11 | 0.60 | 16.00 | 18.50 | 0.56 | 18.50 | 15.99 | 18.84 | 0.56 | 0.54 | 15.09 | 2.79 | good |
| 12 | 0.60 | 15.00 | 17.00 | 0.56 | 15.00 | 15.00 | 17.09 | 0.56 | 0.88 | 15.09 | 2.78 | good |
| 13 | 0.60 | 14.00 | 15.50 | 0.56 | 15.00 | 14.00 | 15.52 | 0.56 | 0.97 | 15.09 | 2.78 | good |
| 14 | 0.60 | 13.00 | 14.00 | 0.56 | 15.00 | 13.00 | 14.01 | 0.56 | 0.99 | 15.09 | 2.78 | good |
| 15 | 0.60 | 12.00 | 12.50 | 0.56 | 15.00 | 12.00 | 12.50 | 0.56 | 1.00 | 15.09 | 2.78 | good |
| 16 | 0.60 | 16.00 | 18.50 | 0.56 | 18.50 | 15.99 | 18.84 | 0.56 | 0.54 | 15.09 | 2.79 | good |
| 17 | 0.60 | 15.00 | 17.00 | 0.56 | 15.00 | 15.00 | 17.09 | 0.56 | 0.88 | 15.09 | 2.78 | good |
| 18 | 0.60 | 14.00 | 15.50 | 0.56 | 15.00 | 14.00 | 15.52 | 0.56 | 0.97 | 15.09 | 2.78 | good |
| 19 | 0.60 | 13.00 | 14.00 | 0.56 | 15.00 | 13.00 | 14.01 | 0.56 | 0.99 | 15.09 | 2.78 | good |
| 20 | 0.60 | 12.00 | 12.50 | 0.56 | 15.00 | 12.00 | 12.50 | 0.56 | 1.00 | 15.09 | 2.78 | good |
| 21 | 0.60 | 16.00 | 18.50 | 1.20 | 15.00 | 15.99 | 18.16 | ** | 2.13 | 15.09 | 2.61 | bad |
| 22 | 0.60 | 15.00 | 17.00 | 1.20 | 15.00 | 14.95 | 16.29 | ** | ** | 15.06 | 2.11 | bad |
| 23 | 0.60 | 14.00 | 15.50 | 1.20 | 15.00 | 14.00 | 15.50 | 1.20 | 0.98 | 15.07 | 2.74 | good |
| 24 | 0.60 | 13.00 | 14.00 | 1.20 | 15.00 | 13.00 | 14.02 | 1.20 | 0.99 | 15.06 | 2.71 | good |
| 25 | 0.60 | 12.00 | 12.50 | 1.20 | 15.00 | 12.00 | 12.51 | 1.20 | 1.00 | 15.05 | 2.69 | good |
| 26 | 0.60 | 16.00 | 18.50 | 2.30 | 15.00 | 16.00 | ** | ** | 0.08 | 15.06 | 2.71 | bad |
| 27 | 0.60 | 15.00 | 17.00 | 2.30 | 15.00 | 15.00 | 15.43 | 3.42 | 0.99 | 15.92 | 3.87 | bad |
| 28 | 0.60 | 14.00 | 15.50 | 2.30 | 15.00 | 14.00 | 15.01 | 3.03 | 1.12 | 16.30 | 5.15 | bad |
| 29 | 0.60 | 13.00 | 14.00 | 2.30 | 15.00 | 13.00 | 13.68 | 2.65 | 1.16 | ** | ** | good |
| 30 | 0.60 | 12.00 | 12.50 | 2.30 | 15.00 | 12.00 | 12.41 | 2.39 | 1.06 | ** | ** | good |
| 31 | 0.60 | 16.00 | 18.50 | 3.40 | 15.00 | 15.99 | ** | ** | ** | 15.10 | 2.58 | bad |
| 32 | 0.60 | 15.00 | 17.00 | 3.40 | 15.00 | 14.98 | ** | ** | ** | 15.10 | 2.13 | bad |
| 33 | 0.60 | 14.00 | 15.50 | 3.40 | 15.00 | 14.00 | 15.26 | 4.95 | 0.67 | 15.38 | 0.69 | bad |
| 34 | 0.60 | 13.00 | 14.00 | 3.40 | 15.00 | 13.00 | 15.26 | 4.96 | 0.67 | 13.97 | 0.65 | bad |
| 35 | 0.60 | 12.00 | 12.50 | 3.40 | 15.00 | 12.00 | 12.42 | 2.07 | 1.19 | ** | ** | good |
| 36 | 0.60 | 16.00 | 18.50 | 4.50 | 15.00 | 16.00 | ** | ** | ** | 15.03 | 2.81 | bad |

Table A1 - continued

|  | Seeing (arcsec) | Simulation parameters |  |  | $m_{\text {disc }}$ | $m_{\text {nuclear }}$ | $m_{\text {bulge }}$ | Retrieved parameters |  |  | $\begin{gathered} R_{\mathrm{s}} \\ (\operatorname{arcsec}) \end{gathered}$ | Model quality |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | $m_{\text {nuclear }}$ | $m_{\text {bulge }}$ | $R$ e (arcsec) |  |  |  | $R_{\text {e }}$ <br> (arcsec) | $n$ | $m_{\text {disc }}$ |  |  |
| 37 | 0.60 | 15.00 | 17.00 | 4.50 | 15.00 | 15.00 | ** | ** | ** | 14.91 | 2.79 | bad |
| 38 | 0.60 | 14.00 | 15.50 | 4.50 | 15.00 | 14.00 | 15.12 | 3.91 | 0.86 | 15.43 | 3.54 | good |
| 39 | 0.60 | 13.00 | 14.00 | 4.50 | 15.00 | 13.00 | 13.94 | 4.13 | 0.93 | 15.25 | 3.92 | good |
| 40 | 0.60 | 12.00 | 12.50 | 4.50 | 15.00 | 12.00 | 12.59 | 4.62 | 0.99 | 14.45 | 2.34 | good |
| 41 | 0.70 | 16.00 | 18.50 | 1.20 | 15.00 | 15.99 | 19.21 | 0.00 | 15.28 | 15.09 | 2.62 | bad |
| 42 | 0.70 | 15.00 | 17.00 | 1.20 | 15.00 | 15.00 | 17.50 | 1.32 | 1.02 | 15.09 | 2.80 | good |
| 43 | 0.70 | 14.00 | 15.50 | 1.20 | 15.00 | 14.02 | 15.04 | 2.42 | 2.37 | 16.33 | 3.82 | good |
| 44 | 0.70 | 13.00 | 14.00 | 1.20 | 15.00 | 13.00 | 14.02 | 1.21 | 0.99 | 15.07 | 2.70 | good |
| 45 | 0.70 | 12.00 | 12.50 | 1.20 | 15.00 | 12.00 | 12.51 | 1.21 | 1.00 | 15.06 | 2.68 | good |
| 46 | 0.80 | 16.00 | 18.50 | 2.30 | 15.00 | 16.00 | 16.43 | 70.00 | 0.04 | 15.10 | 2.62 | bad |
| 47 | 0.80 | 15.00 | 17.00 | 2.30 | 15.00 | 14.99 | 16.51 | 2.72 | 0.92 | 15.25 | 2.95 | bad |
| 48 | 0.80 | 14.00 | 15.50 | 2.30 | 15.00 | 14.00 | 15.02 | 3.02 | 1.11 | 16.31 | 4.94 | bad |
| 49 | 0.80 | 13.00 | 14.00 | 2.30 | 15.00 | 13.00 | 13.68 | 2.65 | 1.16 | 26.10 | 0.01 | good |
| 50 | 0.80 | 12.00 | 12.50 | 2.30 | 15.00 | 12.00 | 12.41 | 2.39 | 1.06 | ** | ** | good |
| 51 | 0.90 | 16.00 | 18.50 | 3.40 | 15.00 | 16.02 | 25.61 | ** | 6.35 | 15.10 | 2.55 | bad |
| 52 | 0.90 | 15.00 | 17.00 | 3.40 | 15.00 | 15.00 | 16.69 | ** | ** | 14.96 | 2.58 | bad |
| 53 | 0.90 | 14.00 | 15.50 | 3.40 | 15.00 | 14.00 | 15.68 | 5.29 | 1.20 | 14.95 | 2.22 | good |
| 54 | 0.90 | 13.00 | 14.00 | 3.40 | 15.00 | 13.00 | 13.78 | 3.46 | 0.97 | 15.99 | 4.15 | good |
| 55 | 0.90 | 12.00 | 12.50 | 3.40 | 15.00 | 12.00 | 12.41 | 3.47 | 1.01 | 17.07 | ** | good |
| 56 | 1.00 | 16.00 | 18.50 | 0.10 | 15.00 | 15.89 | 17.28 | 3.30 | 0.70 | 15.24 | 3.02 | bad |
| 57 | 1.00 | 15.00 | 17.00 | 0.10 | 15.00 | 14.99 | 17.00 | 1.14 | ** | 15.09 | 2.87 | bad |
| 58 | 1.00 | 14.00 | 15.50 | 0.10 | 15.00 | 13.77 | 18.46 | ** | ** | 15.09 | 2.78 | bad |
| 59 | 1.00 | 13.00 | 14.00 | 0.10 | 15.00 | 12.72 | 15.22 | ** | ** | 15.12 | 2.94 | bad |
| 60 | 1.00 | 12.00 | 12.50 | 0.10 | 15.00 | 11.60 | 13.64 | ** | ** | 15.10 | 2.82 | bad |
| 61 | 1.00 | 16.00 | 18.50 | 0.56 | 18.50 | 16.03 | 17.52 | 1.77 | ** | 15.11 | 2.85 | bad |
| 62 | 1.00 | 15.00 | 17.00 | 0.56 | 15.00 | 15.00 | 17.02 | 0.54 | 1.09 | 15.09 | 2.79 | good |
| 63 | 1.00 | 14.00 | 15.50 | 0.56 | 15.00 | 14.00 | 15.51 | 0.55 | 1.00 | 15.09 | 2.78 | good |
| 64 | 1.00 | 13.00 | 14.00 | 0.56 | 15.00 | 13.00 | 14.00 | 0.56 | 1.00 | 15.09 | 2.78 | good |
| 65 | 1.00 | 12.00 | 12.50 | 0.56 | 15.00 | 12.00 | 12.50 | 0.56 | 1.00 | 15.09 | 2.78 | good |
| 66 | 1.00 | 16.00 | 18.50 | 1.20 | 15.00 | 16.21 | 16.87 | ** | ** | 15.08 | 2.63 | bad |
| 67 | 1.00 | 15.00 | 17.00 | 1.20 | 15.00 | 14.98 | 16.24 | ** | 0.99 | 15.06 | 2.25 | bad |
| 68 | 1.00 | 14.00 | 15.50 | 1.20 | 15.00 | 14.00 | 15.57 | 1.20 | 0.94 | 15.07 | 2.72 | good |
| 69 | 1.00 | 13.00 | 14.00 | 1.20 | 15.00 | 13.00 | 14.03 | 1.20 | 0.98 | 15.05 | 2.69 | good |
| 70 | 1.00 | 12.00 | 12.50 | 1.20 | 15.00 | 12.00 | 12.51 | 1.20 | 0.99 | 15.04 | 2.66 | good |
| 71 | 1.00 | 16.00 | 18.50 | 3.40 | 15.00 | 16.00 | ** | ** | ** | 15.07 | 2.63 | bad |
| 72 | 1.00 | 15.00 | 17.00 | 3.40 | 15.00 | 14.99 | 15.99 | 3.52 | 0.80 | 15.43 | 3.17 | bad |
| 73 | 1.00 | 14.00 | 15.50 | 3.40 | 15.00 | 14.00 | 17.12 | ** | 0.12 | 14.62 | 2.29 | bad |
| 74 | 1.00 | 13.00 | 14.00 | 3.40 | 15.00 | 13.01 | 13.67 | 3.54 | 1.03 | 15.32 | ** | good |
| 75 | 1.00 | 12.00 | 12.50 | 3.40 | 15.00 | 12.00 | 12.41 | 3.47 | 1.01 | 16.25 | ** | good |
| 76 | 1.00 | 16.00 | 18.50 | 4.50 | 15.00 | 16.00 | 28.53 | ** | 0.41 | 15.04 | 2.69 | bad |
| 77 | 1.00 | 15.00 | 17.00 | 4.50 | 15.00 | 15.00 | 19.53 | ** | 1.16 | 14.92 | 2.79 | bad |
| 78 | 1.00 | 14.00 | 15.50 | 4.50 | 15.00 | 14.00 | 15.87 | 3.63 | 0.78 | 14.90 | 3.06 | good |
| 79 | 1.00 | 13.00 | 14.00 | 4.50 | 15.00 | 13.00 | 14.20 | 4.84 | 0.98 | 14.69 | 2.43 | good |
| 80 | 1.00 | 12.00 | 12.50 | 4.50 | 15.00 | 12.00 | 12.41 | 4.50 | 0.99 | ** | ** | good |
| 81 | 1.10 | 14.00 | 13.20 | 0.10 | 15.00 | 13.74 | 13.35 | 0.11 | 0.64 | 15.09 | 2.79 | good |
| 82 | 1.10 | 14.00 | 14.00 | 0.10 | 15.00 | 13.55 | 14.62 | ** | ** | 15.11 | 2.94 | bad |
| 83 | 1.10 | 14.00 | 15.00 | 0.10 | 15.00 | 13.68 | 16.95 | 0.12 | ** | 15.11 | 2.88 | bad |
| 84 | 1.10 | 14.00 | 15.50 | 0.10 | 15.00 | 13.79 | 17.57 | 0.50 | ** | 15.09 | 2.78 | bad |
| 85 | 1.10 | 14.00 | 14.50 | 0.10 | 15.00 | 13.59 | 15.91 | 0.13 | 4.11 | 15.10 | 2.79 | bad |
| 86 | 1.10 | 14.00 | 13.20 | 0.40 | 15.00 | 13.95 | 13.23 | 0.40 | 0.95 | 15.09 | 2.75 | good |
| 87 | 1.10 | 14.00 | 14.00 | 0.40 | 15.00 | 13.95 | 14.05 | 0.40 | 0.91 | 15.09 | 2.76 | good |
| 88 | 1.10 | 14.00 | 15.00 | 0.40 | 15.00 | 13.96 | 15.11 | 0.40 | 0.81 | 15.09 | 2.76 | good |
| 89 | 1.10 | 14.00 | 13.00 | 0.40 | 15.00 | 13.94 | 13.02 | 0.40 | 0.96 | 15.09 | 2.75 | good |
| 90 | 1.10 | 14.00 | 14.50 | 0.40 | 15.00 | 13.96 | 14.58 | 0.40 | 0.87 | 15.09 | 2.76 | good |

Notes. Column (1) gives the simulation number; (2) the simulation seeing; (3) and (4) the simulated nuclear and bulge magnitudes; (5) the simulated bulge effective radius; (6) the simulated disc magnitude; (7) the modelled nuclear magnitude; (8) the modelled bulge magnitude; (9) the modelled bulge effective radius; (10) the modelled bulge Sérsic index; (11) the modelled disc magnitude; (12) the modelled disc scale length; (13) the quality with which the simulation was modelled.
The ${ }^{\text {'**' }}$ symbol shows a physically improbable parameter.
All the simulated galaxies have bulges with Sérsic indexes $(n=1)$ and discs with scale lengths ( $R_{\mathrm{S}}=3.0$ arcsec).

## APPENDIX B: MORPHOLOGICAL ANALYSIS

## B1 Models

## B2 Bar test

Figure B1.


Figure B1. Surface brightness profile decomposition of the host galaxies in the sample. The top left panel shows the observed image. The top middle panel shows the image of the model used to describe the surface brightness distribution. The top right panel shows the residual image. Middle: Radial profile of the surface brightness distribution. The symbols and the main parameters of the models are explained in the plots. For the complete Appendix see the online material. Bottom: Residuals.


Figure B2. Radial variation of ellipticity $\epsilon$ (blue circles) and of PA (red squares) derived by ellipse fitting to the galaxy isophotes. Only the host galaxies that fulfil the criteria to identify bars are shown (see Section 4.1). For the complete Appendix see the online material.

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[^1]:    ${ }^{1}$ http://www.not.iac.es/instruments/notcam

[^2]:    ${ }^{2}$ we consider a star bright if it is brighter than the target to fit. A PSF model made from a bright star can fit the wings of the target (and beyond) and its nucleus. Otherwise, a faint star (fainter that the target) will only be able to fit the nucleus and maybe part of the wings.

[^3]:    ${ }^{3}$ In contrast to this work, the original Kormendy relation uses the average surface brightness within the bulge effective radius.

